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SOME USEFUL GENERALIZATIONS IN

MARKOV RENEWAL PROCESSES

by

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DEPARTMENT OF STATISTICS  
Southern Methodist University

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MARKOV RENEWAL PROCESSES

Approved by:

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SOME USEFUL GENERALIZATIONS IN

MARKOV RENEWAL PROCESSES

A Dissertation Presented to the Faculty of the Graduate School

of

Southern Methodist University

in

Partial Fulfillment of the Requirements

for the degree of

Doctor of Philosophy

with a

Major in Statistics

by

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(B.A., University of Dallas, 1964)

(M.S., Southern Methodist University, 1968)

September 22, 1969

Consider a particular system which at time  $t$  may be in one of a finite number of distinguishable states, labeled for convenience by  $1, 2, \dots, m$ . Once the system enters a particular state, say  $i$ , it instantaneously selects the next state to be visited, say  $j$ , with probability  $p_{ij}$ . However, transition to state  $j$  occurs after holding in state  $i$  for a random time (sojourn time) whose distribution function is  $F_{ij}(\cdot)$ . These processes are known as Semi-Markov Processes and the associated renewal process is called a Markov Renewal Process. In this paper we introduce a new counting process which at time  $t$  counts the number of times the system has made a one-step transition from state  $i$  to state  $j$ ,  $i, j = 1, 2, \dots, m$ . The matrix  $N(t)$  denotes these counts. The distribution and moments of  $N(t)$  are derived and cumulative processes associated with  $N(t)$  are discussed. We also extend these results to arbitrary intervals of the form  $(t_0, t_0 + t]$ . Certain limiting results are given and some important special cases discussed. Several open problems are given in the summary chapter.

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Adviser: Associate Professor A. M. Kshirsagar

Some Useful Generalizations in Markov Renewal Processes

B.A., University of Dallas,  
 1964  
 M.S., Southern Methodist  
 University, 1968

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one may loosely describe a S.M.P. as a Markov chain in which the time axis  
 Note that a S.M.P. is Markovian only at transition instants and thus

order Markov Renewal Processes.

and Kshirsagar and Gupta (1968) have studied these processes called zero-  
 are of zero-order arises when  $F_{ij}^{(n)}$  depends either on  $i$  or  $j$ . Pyke (1962)

visited. A special class of M.R.P.'s in which the imbedded Markov chains

both the current state of the system as well as the next state to be

$F_{ij}^{(n)}$ . Note that, in general, the sojourn time distribution depends on

state  $i$  for a random time (sojourn time) whose distribution function is

ability  $p_{ij}$ . However, transition to state  $j$  occurs after holding in

instantaneously selects the next state to be visited, say  $j$ , with prob-

1, 2, ...,  $m$ . Once the system enters a particular state, say  $i$ , it

of a finite number of distinguishable states, labeled for convenience by

In general, consider a particular system which at time  $t$  may be in one

Examples of M.R.P.'s abound in the physical and social sciences.

by Pyke and Moore (1968).

related problem of estimation is treated at some length in a recent paper

is due to Pyke (1961a, 1961b) and Pyke and Schaufele (1964). The asso-

S.M.P.'s as well as the introduction of Markov Renewal Processes (M.R.P.)

pendently by Levy (1954) and Smith (1955). The first thorough study of

The theory of Semi-Markov Processes (S.M.P.) was introduced inde-

## 1. Introduction

### PRELIMINARIES

#### CHAPTER I

Let  $\{X_n\}$  be a Markov chain with finite state space denoted by  $\{e_i; i = 1, 2, \dots, m\}$ , where  $e_j$  is the  $m \times 1$  vector of zeroes except for the  $j$ th entry which is one. Thus  $\bar{e}_j = e_j$  means that at the  $n$ th transition the system is in state  $j$ . Denote the  $m \times 1$  vector of initial probabilities,  $(a_1, a_2, \dots, a_m)$ , by  $\bar{e}_0 = e_0$ ,  $i = 1, 2, \dots, m$ .

$$\text{Pr}\{X_n = e_j, X_{n-1} < x | X_0 = e_i\} = P_{ij}^{(n)}(x), \quad \text{where } P_{ij}^{(n)}(x) \text{ is the } n\text{-th power of the transition matrix } P_{ij}(x).$$

$$\text{Pr}\{X_n = e_j, X_{n-1} < x | X_0 = e_i, X_{n-1} = e_k, \dots, X_1 = e_l\} = P_{ijk\dots l}^{(n)}(x),$$

Let  $\{X_n\}$  be a sequence of non-negative random variables such that

$$\text{Pr}\{X_0 = e_i\} = a_i, \quad i = 1, 2, \dots, m. \quad (1.2.1)$$

Let  $\{X_n; n > 0\}$  be a Markov chain with finite state space denoted by  $\{e_i; i = 1, 2, \dots, m\}$ , where  $e_j$  is the  $m \times 1$  vector of zeroes except for the  $j$ th entry which is one. Thus  $\bar{e}_j = e_j$  means that at the  $n$ th transition the system is in state  $j$ . Denote the  $m \times 1$  vector of initial probabilities,  $(a_1, a_2, \dots, a_m)$ , by  $\bar{e}_0 = e_0$ ,  $i = 1, 2, \dots, m$ .

## 2. Definitions and Notation

chapters.

We present the known results in M.R.P.'s that will be of use in later chapters. In this chapter M.R.P. with finite state space as defined by Pyke (1961b). In this chapter which may arise in the study of S.M.P.'s. Our point of departure is the following chapters our main interest centers on distributions

line simply by letting  $m = 1$ .

Also note that M.R.P. theory embodies ordinary renewal theory on the real line. The S.M.P. reduces to an ordinary Markov chain with finite state space. Then the S.M.P. reduces to a Markov process. If the  $F_{ij}^{(n)}(\cdot)$  are degenerate, has been randomly transformed. If the  $F_{ij}^{(n)}(\cdot)$  are exponential distributions,



The following notation will be used throughout the paper. Random variables and their distribution functions will be denoted by upper case letters and their corresponding Laplace-Stieltjes Transforms (L.S.T) by

$\{m, \bar{a}, \bar{Q}\}$ , where  $\bar{Q}$  is the  $m \times m$  matrix of the  $\bar{Q}_{ij}(\cdot)$ .

shall say that the M.R.P.  $\{\bar{N}(t); t \geq 0\}$  is determined by the triplet the number of times the system enters state  $j$  in the interval  $(0, t]$ . We Then the process  $\{\bar{N}(t); t \geq 0\}$  is a M.R.P. whose  $j^{\text{th}}$  component represents

(1.2.5)

$$\bar{N}(t) = \begin{pmatrix} N_1(t) \\ N_2(t) \\ \dots \\ N_m(t) \\ 0 \end{pmatrix} = \begin{pmatrix} N_1(t) \\ N_2(t) \\ \dots \\ N_m(t) \\ 0 \end{pmatrix} + \dots + \begin{pmatrix} N_1(t) \\ N_2(t) \\ \dots \\ N_m(t) \\ 0 \end{pmatrix}$$

$(0, t]$ . Define the vector-valued random variable

Then  $N(t)$  is the number of transitions that the system makes in the interval

(1.2.4)  $N(t) = \max\{n; S_n \leq t\}$ .

$\{\bar{J}(t); t \geq 0\}$ . Let

$S_n \leq t < S_{n+1}$ .  $X_n$  is called the  $n^{\text{th}}$  sojourn time of the process

$\{\bar{J}^{-n}; X_n; n \geq 0\}$  is equivalent to the S.M.P.  $\{\bar{J}(t); t \geq 0\} = \bar{J}^{-n}$  for

$n \geq 0, X_0 = 0$ , be the time of the  $n^{\text{th}}$  transition. Then the process

$H_i^{\bar{J}}(+\infty) = 1, i = 1, 2, \dots, m$ . Further, let  $S_n = X_0 + X_1 + \dots + X_n$ ,

aperiodic class, which will be assumed throughout this paper, then

$\bar{Q}_{ij}^{\bar{J}}(+\infty) \leq 1$ . If the imbedded Markov chain is an irreducible recurrent

The probability of ultimate transition into state  $j$  is given by

(1.2.3)  $H_i^{\bar{J}}(x) = \sum_{j=1}^m \bar{Q}_{ij}^{\bar{J}}(x)$ .

interval  $(0, x]$  is given by

can also write that the probability of transition to any state in the

lower case letters; i.e., the distribution functions  $Q_{1j}(t)$  have L.S.T.'s

given by

$$Q_{1j}(s) = \int_0^\infty e^{-st} dQ_{1j}(t), \quad (1.2.6)$$

whenever they exist. Also  $q(s)$  will denote the  $m \times m$  matrix of the  $q_{1j}(s)$ .

Let  $Q$  and  $R$  be matrix-valued functions and define a convolution operation

$Q * R$  by

$$Q * R = \left( \sum_{k=1}^m Q_{1k} * R_{kj} \right),$$

$$Q_{1k} * R_{kj} = \int_0^t Q_{1k}(t-u) dR_{kj}(u), \quad (1.2.7)$$

which is the same as matrix multiplication except the product is replaced

with convolution. Let  $E$ ,  $\bar{e}_{1j}$ , and  $\bar{e}_j$  denote respectively an  $m \times m$  matrix

of ones, an  $m \times 1$  vector of ones, an  $m \times m$  null matrix except for the

$(1, j)$ th element which equals one, and an  $m \times 1$  vector of zeroes whose  $j$ th

component equals one. If  $B$  is an  $m \times m$  matrix, then  $B^0$  and  $B^d$  are,

respectively,

$$B^d = (\delta_{1j}^d)_{1j}, \quad B^0 = B - B^d. \quad (1.2.8)$$

Pyke (1961b) derives the distribution and first moment of  $N_j(t)$  in

terms of generating functions. We extend his results and derive the dis-

tribution and first two moments of  $\bar{N}(t)$  in terms of generating functions.

Using results from Kshirsagar and Gupta (1967) we find asymptotic expressions

for the moments of  $\bar{N}(t)$  from series expansions of the L.S.T.'s of the moments.

The resulting expressions are in terms of the basic quantities specified

by  $(m, \bar{a}, \bar{Q})$  rather than in terms of mean recurrence distributions as

$$(1.3.2) \quad = \mathbb{P}(x) \mathbb{P}(x+1) - \mathbb{P}(x) \mathbb{P}(x+1) \cdot$$

$$= 1 - \mathbb{P}(x+1) \mathbb{P}(x) \{1 - \mathbb{P}(x)\}$$

$$\Pr\{N(t) = r\} = \Pr\{N(t) > r + 1\} - \Pr\{N(t) > r\}$$

where  $\mathbb{P}(x)$  is the  $r$ -fold convolution of  $F(t)$ . From (1.3.1) we obtain

$$(1.3.1) \quad \Pr\{N(t) > r\} = \Pr\{S_r > t\} = 1 - \mathbb{P}(x)$$

$(0, t]$  and  $S_r$  denotes the time of the  $r$ th renewal, then

distribution function  $F(t)$ . If  $N(t)$  denotes the number of renewals in times between renewals are independent and identically distributed with a one unit system. For an ordinary renewal process we assume that the

Let  $F(t)$  be the distribution function of the time between renewals in

of the following discussion is given by Cox (1962).

brief account of Renewal Theory in one dimension. A more detailed account

real line. So that these similarities may be pointed out we present a

A close analogy exists between M.R.P.'s and Renewal Processes on the

### 3. Renewal Theory

cases are discussed.

arbitrary. Certain limiting results are given and some important special

for the interval  $(0, t]$  are extended to the interval  $(t_0, t_0 + t]$ , for  $t_0$

processes associated with  $\bar{N}(t)$  and  $N(t)$  are discussed. Lastly, the results

terms of I.S.T.'s and asymptotic means and variances derived. Cumulative

matrix of the  $N_{ij}^1(t)$ . The distribution and moments of  $N(t)$  are derived in

from state  $i$  to state  $j$  in the interval  $(0, t]$  and let  $N(t)$  be the  $m \times m$

Let  $N_{ij}^1(t)$  be the number of times the system makes a one step transition

Pyke (1961b) has done. We also discuss a generalization of  $\bar{N}(t)$ ; namely,

The statistical properties of  $N(t)$  are often more easily discussed from the probability generating function (p.g.f.) of  $N(t)$  and so we define

$$G(t, z) = \sum_{r=0}^{\infty} z^r \Pr\{N(t) = r\}$$

$$= \sum_{r=0}^{\infty} z^r \Pr(x)(t) - \sum_{r=0}^{\infty} z^r \Pr(r+1)(t)$$

$$= 1 + \sum_{r=1}^{\infty} z^r \Pr(x)(t) - \sum_{r=1}^{\infty} z^{r-1} \Pr(x)(t)$$

$$(1.3.3) \quad = 1 + \sum_{r=1}^{\infty} z^{r-1} \Pr(x)(t) - \sum_{r=1}^{\infty} z^{r-1} \Pr(x)(t)$$

Define

$$F(s) = \int_0^{\infty} e^{-st} d_q F(t)$$

as the L.S.T. of  $F(t)$ . The reader will note that Cox (1962) defines

$$(1.3.5) \quad F^*(s) = \int_0^{\infty} e^{-st} F(t) dt$$

as the L.S.T. of  $f(t)$ . However in Cox's notation we have that the L.S.T. of  $F(t)$  is

$$(1.3.6) \quad \int_0^{\infty} e^{-st} F(t) dt = \int_0^{\infty} \frac{1}{s} e^{-st} F(t) dt = \frac{1}{s} F^*(s)$$

and so we make the substitution

$$(1.3.7) \quad F(s) = \frac{1}{s} F^*(s)$$

In order that our notation agrees with that of Cox.

function. Now for an arbitrary function  $k(x)$ ,  
 forced to consider the asymptotic behavior of the transform of the renewal  
 in all but the simplest cases  $h_0(s)$  cannot be inverted exactly and one is

$$(1.3.14) \quad h^m(s) = \frac{\partial g^m(s, z)}{\partial z} \Big|_{z=1} = f^m_1(s) [1 - f(s)]^{-1}$$

$$(1.3.13) \quad h_0(s) = \frac{\partial g_0(s, z)}{\partial z} \Big|_{z=1} = f(s) [1 - f(s)]^{-1}$$

The transform of the renewal function,  $H(t) = E(N(t))$ , is in our notation

$$(1.3.12) \quad g^m(s, z) = 1 - (1 - z) f^m_1(s) [1 - zf(s)]^{-1}$$

and

$$(1.3.11) \quad g_0(s, z) = [1 - f(s)] [1 - zf(s)]^{-1}$$

In our notation

$$(1.3.10) \quad g^m(s, z) = \frac{[1 - zf^*(s)]}{1 - zf^*(s) + zf^*_1(s)}$$

obtain

for the ordinary renewal process. For the modified renewal process we

$$(1.3.9) \quad = \frac{[1 - zf^*(s)]}{1 - f^*(s)}$$

$$= \frac{s}{1} [1 - f^*(s)] \left( [1 + zf^*(s)] + [zf^*(s)]^2 + \dots \right)$$

$$(1.3.8) \quad g_0(s, z) = \frac{s}{1} + \sum_{r=1}^{\infty} z^{-r} (z - 1) f^*(s)^r$$

Taking the L.S.T. of (1.3.3) we have in Cox's notation

$$V_{1j}^j(k;t) = \Pr\{N_j(t) = k | \bar{J}(0) = e^{-1}\} = \Pr\{N_j(t) = k\} \quad (1.4.1)$$

Define

1961.

We turn now to a discussion of the main results in M.R.P.'s that have appeared in the literature since the publication of Pyke's papers in

#### 4. Basic Results in M.R.P.'s

with ordinary renewal theory.

As M.R.P. theory is developed we shall point out some of the analogies. At this point we require no further results from ordinary renewal theory.

$$H_0(t) = \frac{\mu}{t} + \frac{\sigma^2}{2} \frac{2\mu}{t^2} + o(1) \quad (1.3.17)$$

whose inverse is, for large  $t$ ,

$$h_0(s) = \frac{1}{s\mu} + \frac{\sigma^2}{2} \frac{2\mu}{s^2} + o(1) \quad (1.3.16)$$

Substituting into (1.3.13) and simplifying gives

$$f(s) = 1 - s\mu + \frac{1}{2} s^2 (\mu^2 + \sigma^2) + o(s^2) \quad (1.3.15)$$

we have for large  $t$ ; i.e., small  $s$ ,

of  $h_0(s)$  for small  $s$ . Letting  $\mu$  and  $\sigma^2$  be the mean and variance of  $F(t)$  the function  $H_0(t)$  for large  $t$  can be studied by examining the behavior if  $x$  is large,  $k(s)$  is negligible unless  $s$  is small. Hence the form of

$$k(s) = \int_0^\infty e^{-sx} d_x k(x) \quad .$$

$$\begin{aligned}
&= 1 - H_{1j}^T(t) + (z - 1) \delta_{1j}^T(t) * \phi_{1j}^T(z;t) + \sum_{m=1}^{\infty} \delta_{1j}^T(t) * \phi_{1j}^T(z;t) \\
&= 1 - H_{1j}^T(t) + z \delta_{1j}^T(t) * \phi_{1j}^T(z;t) + \sum_{k=0}^{\infty} z^k \delta_{1j}^T(t) * \phi_{1j}^T(z;t) \\
&\quad + z \sum_{k=1}^{\infty} z^{k-1} \delta_{1j}^T(t) * \phi_{1j}^T(z;t) \\
&\quad + z \sum_{k=1}^{\infty} z^{k-1} \delta_{1j}^T(t) * \phi_{1j}^T(z;t) \\
&= 1 - H_{1j}^T(t) + \sum_{r \neq j} \delta_{1j}^T(t) * \phi_{1j}^T(z;t) \\
&\quad + \sum_{k=1}^{\infty} z^{k-1} \delta_{1j}^T(t) * \phi_{1j}^T(z;t)
\end{aligned}$$

(1.4.2) gives

as the L.S.T. of  $\phi_{1j}^T(z;t)$ . Applying (1.4.3) to the system of equations

$$\psi_{1j}^T(z;s) = \int_0^{\infty} e^{-st} d\phi_{1j}^T(z;t) \tag{1.4.4}$$

be the probability generating function (p.g.f.) of the  $V_{1j}^T(k;t)$  and define

$$\phi_{1j}^T(z;t) = \sum_{k=0}^{\infty} z^k V_{1j}^T(k;t) \tag{1.4.3}$$

Let, for  $|z| < 1$ ,  $i, j = 1, 2, \dots, m$ ,

$$\begin{aligned}
V_{1j}^T(k;t) &= \delta_{1j}^T(t) * V_{1j}^T(k-1;t) + \sum_{r \neq j} \delta_{1j}^T(t) * V_{1j}^T(k;t) \\
V_{1j}^T(0;t) &= 1 - H_{1j}^T(t) + \sum_{r \neq j} \delta_{1j}^T(t) * V_{1j}^T(0;t)
\end{aligned} \tag{1.4.2}$$

Then we immediately have that

whose L.S.T. is

$$\psi_{1j}^T(z;s) = I - h^T(s) + (z - 1)q_{1j}^T(s)\phi_{1j}^T(z;s) + \sum_{r=1}^m q_{1r}^T(s)\phi_{1r}^T(z;s) \quad (1.4.5)$$

or in matrix notation

$$\psi^z = (I - q(s))E + (z - 1)q(s)\psi^z + q(s)\psi^z \quad (1.4.6)$$

$(I - q(s))$  is nonsingular hence

$$\psi^z + (I - q(s))\psi^z = E \quad (1.4.7)$$

Taking diagonal elements on both sides of (1.4.7) gives

$$\psi^z + (I - q(s))\psi^z = E$$

so that

$$\psi^z = \left[ I + (I - q(s))\psi^z \right]^{-1} E \quad (1.4.8)$$

Substituting into (1.4.7) gives

$$\psi^z = E - (I - q(s))\psi^z + (I - q(s))\psi^z \quad (1.4.9)$$

Lemma:  $(I - q(s))\psi^z = (I - q(s))\psi^z$

PROOF:  $(I - q(s))\psi^z = (I - q(s))\psi^z$

$$= \psi^z + \dots$$

$$= \psi^z + \dots = (I - q(s))\psi^z$$



$$\begin{aligned} & \cdot \left[ \left\{ \left( (s) \mathbb{B} - \mathbb{I} \right) \right\}^{\mathbb{P}} (z - \mathbb{I}) + \mathbb{I} z \right] \left[ \left\{ \left( (s) \mathbb{B} - \mathbb{I} \right) \right\}^{\mathbb{P}} - \mathbb{I} \right] \\ & x_{-1} \left[ \left\{ \left( (s) \mathbb{B} - \mathbb{I} \right) \right\}^{\mathbb{P}} (z - \mathbb{I}) + \mathbb{I} z \right] (s) \mathbb{B}_{-1} \left( (s) \mathbb{B} - \mathbb{I} \right) (z - \mathbb{I}) + \\ & \quad \left[ \left\{ \left( (s) \mathbb{B} - \mathbb{I} \right) \right\}^{\mathbb{P}} (z - \mathbb{I}) + \mathbb{I} z \right] (s) \mathbb{B}_{-1} \left( (s) \mathbb{B} - \mathbb{I} \right) = z \psi \frac{dz}{z} \end{aligned}$$

Differentiating (1.4.10) gives

$$m(s) = \left. z \frac{dz}{z} \psi \right|_{z=1}$$

$m(s) = (M_{1j}^{1j}(s))$  we have

and let  $m_{1j}^{1j}(s)$  and  $r_{1j}^{1j}(s)$  denote their L.S.T.'s respectively. Setting

$$R_{1j}^{1j}(t) = E\{N_{1j}^{1j}(t) [N_{1j}^{1j}(t) - 1]\} \quad (1.4.12)$$

$$M_{1j}^{1j}(t) = E\{N_{1j}^{1j}(t)\} \quad (1.4.11)$$

define

$\psi_{1j}^{1j}(k;t)$  are easily obtained by differentiation of (1.4.10). In particular

a result obtained by Pyke (1961b). The L.S.T.'s of the moments of the

$$= E - (1 - z) \left( (s) \mathbb{B} - \mathbb{I} \right) \left[ z \mathbb{I} + (1 - z) \left\{ \left( (s) \mathbb{B} - \mathbb{I} \right) \right\}^{\mathbb{P}} \right]_{-1} \quad (1.4.10)$$

$$\psi^z = E - (1 - z) \left( (s) \mathbb{B} - \mathbb{I} \right) \left[ \mathbb{I} + (1 - z) \left\{ \left( (s) \mathbb{B} - \mathbb{I} \right) \right\}^{\mathbb{P}} \right]_{-1} \left[ \mathbb{I} \right]_{-1}$$

Using the Lemma in (1.4.9) gives

$$= (\mathbb{I} - A)_{-1} - \mathbb{I}$$

$$A + A^2 + \dots = (\mathbb{I} + A + A^2 + \dots) - \mathbb{I}$$

Also

Setting  $z = I$  gives

$$(1.4.13) \quad m(s) = (I - q(s))^{-1} q(s)$$

as first derived by Pyke (1961b). In ordinary renewal theory we have

shown that the L.S.T. of  $E\{N(t)\}$  is given by

$$h^0(s) = f(s) [I - f(s)]^{-1}$$

which is the same as  $m(s)$  with the matrix  $q(s)$  in place of the scalar

$f(s)$ .

Note from (1.4.13) that  $(m, \bar{a}, M(t))$  completely determines the M.R.P.

Differentiating a second time and setting  $z = I$  gives the matrix

$$r(s) = (r_{1j}^T(s)) : i.e.,$$

$$r(s) = 2 \left[ (I - q(s))^{-1} q(s) \right]^p \left\{ (I - q(s))^{-1} - I \right\}$$

$$(1.4.14) \quad = 2m(s)^p \{m(s)\} \cdot$$

From (1.4.13) and (1.4.14) the second moments and variance are obtainable.

### 5. Probability Distribution of $\bar{N}(t)$

Cinlar (1968) and Kshirsagar and Gupta (1969a) generalized the results

of 1.4 and derived the joint distribution of the vector  $\bar{N}(t)$  and  $\bar{J}(t)$ . For the

purposes of this paper we derive, in a manner analogous to Cinlar's, the

joint distribution of the vector  $\bar{N}(t)$ . Letting  $\bar{k}$  be a vector of non-

negative integers define

$$(1.5.1) \quad V_{\bar{k};t}^T = \Pr\{\bar{N}(t) = \bar{k} | \bar{J}(0) = \bar{e}_1\} = \Pr\{\bar{N}(t) = \bar{k}\}$$

Let

$$\phi_{\bar{k};t}^T(z_1, z_2, \dots, z_m, z) = \sum_{k_1=0}^{\infty} \dots \sum_{k_m=0}^{\infty} z_1^{k_1} z_2^{k_2} \dots z_m^{k_m} V_{\bar{k};t}^T$$

$$(1.5.1) \quad \bar{e}^{-1} (z(s) \bar{b} - I)_{-1}^{-1} (z(s) \bar{b} - I) = z \bar{\phi}$$

obtaining

that  $(I - q(s)Z)$  is also nonsingular so that we can solve (1.5.4) for  $\bar{\phi}^z$

Note that  $|z^i| \leq 1, i = 1, 2, \dots, m$  and  $(I - q(s))$  nonsingular imply

$$(1.5.4) \quad \bar{\phi}^z = q(s)Z \bar{\phi}^z + (I - q(s)) \bar{e}^{-1}$$

or in matrix notation

$$(1.5.3) \quad \begin{pmatrix} \phi^1(z) \\ \vdots \\ \phi^m(z) \end{pmatrix} = \begin{pmatrix} z \\ \vdots \\ z \end{pmatrix} \begin{pmatrix} \phi^1(z) \\ \vdots \\ \phi^m(z) \end{pmatrix} + \begin{pmatrix} 1 \\ \vdots \\ 1 \end{pmatrix} \bar{e}^{-1} + \sum_{w=1}^{r-1} \begin{pmatrix} \phi^1(z) \\ \vdots \\ \phi^m(z) \end{pmatrix} \bar{e}^{-1} z^w$$

whose L.S.T. is

$$\begin{pmatrix} \phi^1(z) \\ \vdots \\ \phi^m(z) \end{pmatrix} = \begin{pmatrix} z \\ \vdots \\ z \end{pmatrix} \begin{pmatrix} \phi^1(z) \\ \vdots \\ \phi^m(z) \end{pmatrix} + \begin{pmatrix} 1 \\ \vdots \\ 1 \end{pmatrix} \bar{e}^{-1} + \sum_{w=1}^{r-1} \begin{pmatrix} \phi^1(z) \\ \vdots \\ \phi^m(z) \end{pmatrix} \bar{e}^{-1} z^w$$

$\bar{e}^{-1}$  having  $k_r < 0$ . Taking the p.g.f. in (1.5.2) one obtains

where it is understood that summation in (1.5.2) extends over those states

$$(1.5.2) \quad \Delta^T(k;t) = \sum_{w=1}^{r-1} \bar{e}^{-1} \Delta^T(k;t) * \Delta^T(k;t) - \bar{e}^{-1} \Delta^T(k;t)$$

$$\Delta^T(0;t) = 1 - H^T(t)$$

denote  $\bar{\phi}^z = \begin{pmatrix} \phi^1(z) \\ \vdots \\ \phi^m(z) \end{pmatrix}$ . Then

where  $|z^i| \leq 1, i = 1, 2, \dots, m$ . Let  $Z = \text{diag}(z^1, \dots, z^m)$  and

$$\int_{-\infty}^0 e^{-st} \bar{\phi}^z d_t \bar{\phi}^z = \begin{pmatrix} \phi^1(z) \\ \vdots \\ \phi^m(z) \end{pmatrix}$$

and

$$(1.5.9) \quad \bar{c}_{jk}^{-1}(s) = m(s)e_{jk}^{-1} + m(s)e_{jk}^{-1} + m(s)e_{jk}^{-1} \bar{e}^{-1}$$

For  $Z = I$  we obtain the vector  $\bar{c}_{jk}^{-1}(s) = \bar{c}_{jk}^{-1}(s)$  : I.e.,

$$\bar{e}^{-1}((s) \bar{b} - I)_{-1}^{-1} (Z(s) \bar{b} - I) e_{jk}^{-1} (Z(s) \bar{b} - I)_{-1}^{-1} (Z(s) \bar{b} - I) +$$

$$\bar{e}^{-1}((s) \bar{b} - I)_{-1}^{-1} (Z(s) \bar{b} - I) e_{jk}^{-1} (Z(s) \bar{b} - I)_{-1}^{-1} (Z(s) \bar{b} - I) = \frac{z e_{jk}^{-1} z}{z \bar{\psi} z}$$

and

$$\bar{e}^{-1}((s) \bar{b} - I)_{-1}^{-1} (Z(s) \bar{b} - I) e_{jk}^{-1} (Z(s) \bar{b} - I)_{-1}^{-1} (Z(s) \bar{b} - I) = \frac{z e_{jk}^{-1} z}{z \bar{\psi} z}$$

To obtain the product moments (1.5.7) we take the mixed partial  $\partial/\partial z_{jk}^{-1}$ , and evaluate the resulting expression at  $Z = I$ . In particular,

$$(1.5.8) \quad \bar{m}(s) = m(s) \bar{e}^{-1}$$

$$\bar{m}(s) = \bar{m}(s) (I - q(s))_{-1}^{-1} \bar{e}^{-1}$$

Their L.S.T.'s will be denoted by  $\bar{m}(s)$  and  $\bar{c}_{jk}^{-1}(s)$  respectively. Differentiating (1.5.5) and setting  $Z = I$  we obtain, omitting details,

$$(1.5.7) \quad \bar{c}_{jk}^{-1}(t) = E\{N_j^{-1}(t) N_k^{-1}(t) | \bar{J}(0) = \bar{e}^{-1}\} = E\{N_j^{-1}(t) N_k^{-1}(t)\}$$

and

$$(1.5.6) \quad \bar{M}(t) = E\{\bar{N}_j^{-1}(t)\}$$

similar to  $\bar{\psi}_z$  above. Let for ordinary renewal theory is  $(I - f(s)z)_{-1}^{-1} (1 - f(s))$  which is remarkably Note that  $\bar{\psi}_0 = (I - q(s)) \bar{e}^{-1}$  and  $\bar{\psi}_z = \bar{e}^{-1}$  as required. The analogous result

whose  $i$ th element is

$$c_{jk}^{(i)}(s) = m_{ik}^{(j)}(s) + m_{ij}^{(k)}(s) \quad (1.5.10)$$

For  $j = k$ ,

$$c_{jj}^{(i)}(s) = 2m_{ij}^{(j)}(s) + m_{ij}^{(j)}(s) \quad (1.5.11)$$

So that for any  $j$  and  $k$

$$c_{jk}^{(i)}(s) = \delta_{jk}^{(i)} m_{ij}^{(j)}(s) + m_{ik}^{(j)}(s) + m_{ij}^{(k)}(s) \quad (1.5.12)$$

Ginlar (1968) also defines a new vector-valued process which at time  $t$  gives the amounts of times spent in the various states in the interval  $(0, t]$ . He also considers the vector-valued process which gives the amounts of times spent in the various states at the instant of the  $n$ th transition. Various distributions associated with M.R.P.'s in space are introduced and briefly discussed. We do not present details of his results here, since they do not pertain to results in later chapters. It suffices to note that Ginlar's results in  $m$  dimensional space are the exact analogue for a simple renewal process on the real line.

From (1.5.3) or (1.5.5) various marginal distributions are easily obtained. For example, in (1.4.3) let  $Z^* = e + (z - 1)e_j$  to obtain an expression for the marginal distribution of  $N_j(t)$ . Formally we obtain, deleting arguments of the  $\psi_i$ ,

$$\psi_i^{(0)} = 1 - h_i(s) + \sum_{r=1}^m q_{ir}^{(i)}(s) \psi_0^{(r)} + (z - 1) q_{ij}^{(i)}(s) \psi_j^{(0)} \quad (1.5.13)$$

Let  $q^{(i)}(s)$  be the  $i$ th column of the matrix  $q(s)$ . Then (1.5.13) may be written in matrix notation in the form

and  $k$  elements on both sides of (1.5.16) giving the following system where  $\alpha = z_j - 1$ ,  $\beta = z_k - 1$ . As in the previous case we take the  $j$ th

$$(1.5.16) \quad \bar{\psi}_0^j = \alpha \psi_0^j (I - D(s))^{-1} (I - D(s))^{-1} \psi_0^k + \beta \psi_0^k (I - D(s))^{-1} (I - D(s))^{-1} \bar{\psi}_0^k + \bar{e}_j$$

in (1.5.3). Then it is easily shown that the joint marginal distribution of  $(N_j^k(t), N_k^j(t))$  is found by setting  $Z^* = E + (z_j - 1)e_j + (z_k - 1)e_k$  as first derived by Pyke (1961b). Similarly, the joint marginal distribution

$$\psi_0^z = (z - 1)(I - D(s))^{-1} (I - D(s))^{-1} \psi_0^z + (z - 1)(I - D(s))^{-1} (I - D(s))^{-1} \psi_0^z + E + (z - 1)e_j + (z - 1)e_k$$

or matrix notation

$$(1.5.15) \quad \psi_0^z = \frac{(z - 1) a_j^z(s)}{(z - 1) a_j^z(s) + 1} \psi_0^z + \bar{e}_j$$

whose  $i$ th element is

$$\bar{\psi}_0^j = \frac{1 - (z - 1) a_j^j(s)}{(z - 1) (I - D(s))^{-1} (I - D(s))^{-1} \bar{\psi}_0^j + 1} \bar{\psi}_0^j + \bar{e}_j$$

Substituting into (1.5.14) gives

$$\psi_0^j = [1 - (z - 1) a_j^j(s)]^{-1} \psi_0^j$$

Let  $a_j^j(s) = (I - D(s))^{-1} (I - D(s))^{-1}$ . Then

$$\psi_0^j = (z - 1) \psi_0^j (I - D(s))^{-1} (I - D(s))^{-1} \psi_0^j + 1$$

The  $j$ th element of  $\bar{\psi}_0^j$  is then

$$(1.5.14) \quad \bar{\psi}_0^j = (z - 1) \psi_0^j (I - D(s))^{-1} (I - D(s))^{-1} \bar{\psi}_0^j + \bar{e}_j$$

In this section we show how the results of Sections 4 and 5 may be extended to arbitrary intervals  $(t_0, t_0 + t]$ . One often encounters this

6. Extension to Arbitrary Intervals

in agreement with (1.4.13). From (1.5.18) or (1.5.5) one can determine the  $c_{jk}^i(s)$ , the L.S.T.'s of  $B\{N_j^i(t)N_j^k(t)\}$ .

$$\left. \frac{\partial \psi_0^i}{\partial \alpha} \right|_{\alpha=\beta=0} = a_{ij}^i \psi_0^i(s) \quad , \quad \left. \frac{\partial \psi_0^i}{\partial \beta} \right|_{\alpha=\beta=0} = a_{jk}^i \psi_0^i(s) \quad ,$$

Note that  $\psi_0^i$  at  $z_j^k = 1$  equals one as required and that

$$\psi_0^i = \frac{a_{ij}^i(s) \left( 1 - \beta a_{jk}^i(s) + \beta a_{jk}^i(s) + \beta a_{jk}^i(s) \right) \left( 1 - \alpha a_{ij}^i(s) \right) \left( 1 - \alpha \beta a_{jk}^i(s) a_{jk}^i(s) \right)}{\left( 1 - \alpha a_{ij}^i(s) + \alpha a_{jk}^i(s) \right) \left( 1 - \alpha a_{ij}^i(s) + \alpha a_{jk}^i(s) \right)} \quad (1.5.18)$$

we may write

$$\psi_0^i = \alpha \psi_0^i a_{ij}^i(s) + \beta \psi_0^i a_{jk}^i(s) + 1 \quad ,$$

Since

$$\psi_0^i = \frac{\left( 1 - \alpha a_{ij}^i(s) \right) \left( 1 - \beta a_{jk}^i(s) \right) \left( 1 - \alpha \beta a_{jk}^i(s) a_{jk}^i(s) \right)}{\left( 1 - \alpha a_{ij}^i(s) - a_{jk}^i(s) \right)}$$

$$\psi_0^j = \frac{\left( 1 - \alpha a_{ij}^i(s) \right) \left( 1 - \beta a_{jk}^i(s) \right) \left( 1 - \alpha \beta a_{jk}^i(s) a_{jk}^i(s) \right)}{\left( 1 - \beta a_{jk}^i(s) - a_{jk}^i(s) \right)}$$

The solution is easily shown to be

$$\psi_0^j = \alpha \psi_0^j a_{ij}^i(s) + \beta \psi_0^j a_{jk}^i(s) + 1 \quad ,$$

$$\psi_0^k = \alpha \psi_0^k a_{ij}^i(s) + \beta \psi_0^k a_{jk}^i(s) + 1 \quad . \quad (1.5.17)$$

of equations

situation in physical applications since seldom do we observe systems from the start of operation. One is more likely to begin observing a system after it has been in operation for a time  $t_0$ , where  $t_0$  may be large.

Following the derivation given in Kshirsagar and Gupta (1969b) the distribution of the vector  $\bar{N}(t_0, t)$ , whose  $j$ th component is the number of times the system visits the  $j$ th state in  $(t_0, t_0 + t]$ , is derived. Expressions for the L.S.T.'s of the first two moments are also given. As in Sections 4 and 5 denote the moments for the interval  $(t_0, t_0 + t]$  by  $M_{ij}^{(j)}(t_0, t)$ ,  $R_{ij}^{(j)}(t_0, t)$ , and  $C_{ij}^{(j)}(t_0, t)$ . Their L.S.T.'s, with respect to  $t$ , will be denoted by  $m_{ij}^{(j)}(s_0, s)$ ,  $r_{ij}^{(j)}(s_0, s)$ , and  $c_{ij}^{(j)}(s_0, s)$  respectively. Also we define their L.S.T.'s, with respect to  $t_0$ , by

$$\int_0^\infty e^{-s_0 t_0} m_{ij}^{(j)}(s_0, s) dt_0 = m_{ij}^{(j)}(s_0, s),$$

$$\int_0^\infty e^{-s_0 t_0} r_{ij}^{(j)}(s_0, s) dt_0 = r_{ij}^{(j)}(s_0, s),$$

$$\int_0^\infty e^{-s_0 t_0} c_{ij}^{(j)}(s_0, s) dt_0 = c_{ij}^{(j)}(s_0, s),$$

respectively. From Cox (1962, pg. 26) we have that the renewal density, which specifies the mean number of renewals in small intervals near  $t$ , is defined by

$$h_{ij}^{(j)}(t) = \lim_{\Delta t \rightarrow 0^+} \frac{1}{\Delta t} E \left( N_{ij}^{(j)}(t, t + \Delta t) \right)$$

$$= \frac{d}{dt} M_{ij}^{(j)}(t).$$

This represents the probability of a transition from  $i$  to  $j$  in the interval  $(t, t + \Delta t)$ . Also its L.S.T. is given by



$$(1.6.2) \quad \int_0^{t_0} h_{1j}(t) \frac{dx}{dt} (n) - \int_0^{t_0} h_{1j}(t) \frac{dx}{dt} (n) + x_0^j (n) du$$

$$\frac{dx}{dt} (n) \frac{dx}{dt} (n) = \left( \frac{dx}{dt} (n) \right) = \frac{dx}{dt} (n) + x_0^j (n)$$

transition after  $t_0$ . Thus we have shown

Summing the contributions from (a) and (b) gives the p.d.f. of the first

$$\int_0^{t_0} h_{1j}(t) \frac{dx}{dt} (n) - \int_0^{t_0} h_{1j}(t) \frac{dx}{dt} (n) + x_0^j (n) du$$

occurs with probability

occurs in the interval  $(n, n + \Delta x_0)$ . This

$(t_0 - n, t_0 - n + \delta)$  and then transition to state  $k$

b) for some  $n$  transition to state  $j$  occurs in the interval

$$\frac{dx}{dt} (n) \frac{dx}{dt} (n) + x_0^j (n), \text{ or}$$

a) the first transition occurs in the interval  $(t_0 + x_0^j, t_0 + x_0^j + \Delta x_0)$ . This occurs with probability

a) the first transition occurs in the interval  $(t_0 + x_0^j, t_0 + x_0^j + \Delta x_0)$ .

pg. 62) that

transition to occur in the interval  $(x_0^j, x_0^j + \Delta x_0)$ , we have from Cox (1962,

to  $k$  with sojourn time, measured from  $t_0$ , less than or equal to  $x_0^j$ . For

the probability that the first transition after time  $t_0$  is from state  $j$

sequent transitions in  $(t_0, t_0 + t]$ . Let  $Q_{jk}^{1j}(t_0, t_0 + t] = e_{1j}^{-1}$  denote

distribution of the first transition after  $t_0$  is different than all sub-

to S.M.P.'s of the distribution of forward recurrence-time. Note that the

as noted by Kshirsagar and Gupta (1969b). We consider now a generalization

$$(1.6.1) \quad = \left\{ I - Q(s) \right\}^{-1} \left\{ -I - \delta_{1j} \right\}$$

$$h_{1j}(s) = \int_0^\infty e^{-st} h_{1j}(t) dt = m_{1j}(s)$$

which is analogous to the p.d.f. of forward recurrence time in ordinary

renewal theory given by Cox (1962, pg. 63). Let  $q_{jk}^*(s_0)$ ,  $s|\bar{j}(0) = e_1^{-1}$  be

the L.S.T. of (1.6.2), i.e.

$$\int_0^\infty e^{-st} \int_0^\infty e^{-sx} \left[ \frac{dx}{dt} \right]_{t_0} q_{jk}^*(s_0, x|\bar{j}(0)) = e_1^{-1} \left[ \frac{dt_0}{dt} \right]$$

$$= \int_0^\infty e^{-st} \int_0^\infty e^{-sx} \frac{dx}{dt_0} q_{jk}^*(t_0 + x) dt_0$$

$$+ \int_0^\infty e^{-st} \int_0^\infty e^{-sx} \int_0^\infty h_{1j}(t_0 - u) \frac{dx}{dt} q_{jk}^*(u+x) du dt_0$$

(1.6.3)

In the first integral let  $v = t_0 + x$  giving

$$\int_0^\infty e^{-st} \int_0^\infty e^{-s(v-t_0)} \frac{dv}{dt_0} q_{jk}^*(v) dt_0 = \int_0^\infty e^{-sv} \int_0^\infty e^{-s(v-t_0)} \frac{dv}{dt_0} dt_0 q_{jk}^*(v)$$

$$= \int_0^\infty e^{-sv} \left[ 1 - e^{-s(v-s_0)} \right] \frac{dv}{dt_0} q_{jk}^*(v)$$

$$= \left[ \frac{s_0}{s} q_{1j}^*(s) - q_{jk}^*(s_0) \right]$$

(1.6.4)

The second integral can be written in the form

$$\int_0^\infty e^{-sx} \int_0^\infty \frac{dx}{dt} q_{jk}^*(u+x) e^{-s(u+v)} h_{1j}(v) dv du$$

Let  $v = u$  giving

$$\int_0^\infty e^{-sx} \int_0^\infty \frac{dx}{dt} q_{jk}^*(u+x) e^{-s(u+v)} h_{1j}(v) dv du$$

$$= \int_0^\infty e^{-s_0 u} \int_0^\infty e^{-sx} \frac{dx}{dt} q_{jk}^*(u+x) du$$

(1.6.8)

$$\begin{aligned} & \left( \frac{s - 0}{s} \right) \left( \frac{0}{s} \right)_{b}^{k_{f_j}} - \left( \frac{0}{s} \right)_{b}^{k_{f_j}} \left\{ \frac{0}{s} \right\} - \frac{0}{s} + \\ & \left( \frac{s - 0}{s} \right) \left( \frac{0}{s} \right)_{z}^{k_{f_j}} \left( \frac{0}{s} \right)_{b}^{k_{f_j}} \left\{ \frac{0}{s} \right\} = \left( \frac{0}{s} \right)_{z}^{k_{f_j}} \end{aligned}$$

Taking L.S.T.'s with respect to  $t_0$  gives

(1.6.7)

$$\left( \frac{0}{s} \right)_{b}^{k_{f_j}} \left( \frac{0}{s} \right)_{b}^{k_{f_j}} + 1 - \left( \frac{0}{s} \right)_{b}^{k_{f_j}} \left( \frac{0}{s} \right)_{b}^{k_{f_j}} = \left( \frac{0}{s} \right)_{b}^{k_{f_j}} \left( \frac{0}{s} \right)_{b}^{k_{f_j}} = \left( \frac{0}{s} \right)_{b}^{k_{f_j}}$$

Then

respect to  $t$  of the p.g.f. of  $\bar{N}(t, t_0)$  given the initial state is 1. where we have used (1.6.1). Define  $v_{f_j}^T(z, t_0, s)$  as the L.S.T. with

$$\begin{aligned} & \left( \frac{0}{s} \right)_{b}^{k_{f_j}} \left( \frac{0}{s} \right)_{b}^{k_{f_j}} = \frac{s - 0}{s} \left\{ \frac{0}{s} \right\} \\ & \left( \frac{0}{s} \right)_{b}^{k_{f_j}} \left( \frac{0}{s} \right)_{b}^{k_{f_j}} = \left( \frac{0}{s} \right)_{b}^{k_{f_j}} \left( \frac{0}{s} \right)_{b}^{k_{f_j}} + 1 - \left( \frac{0}{s} \right)_{b}^{k_{f_j}} \left( \frac{0}{s} \right)_{b}^{k_{f_j}} \end{aligned}$$

Using (1.6.4) and (1.6.5) in (1.6.3) gives

$$\begin{aligned} & \left( \frac{0}{s} \right)_{b}^{k_{f_j}} \left( \frac{0}{s} \right)_{b}^{k_{f_j}} = \frac{s - 0}{s} \left( \frac{0}{s} \right)_{b}^{k_{f_j}} \\ & \left( \frac{0}{s} \right)_{b}^{k_{f_j}} \left( \frac{0}{s} \right)_{b}^{k_{f_j}} = \int_0^{\infty} \frac{s - 0}{s} e^{-sv} \left[ 1 - e^{-v(s-0)} \right] \frac{d}{dv} \left( \frac{0}{s} \right)_{b}^{k_{f_j}}(v) \\ & \left( \frac{0}{s} \right)_{b}^{k_{f_j}} \left( \frac{0}{s} \right)_{b}^{k_{f_j}} = \int_0^{\infty} e^{-sv} \frac{d}{dv} \left( \frac{0}{s} \right)_{b}^{k_{f_j}}(v) e^{-n_0 + n_0 s} dv \end{aligned}$$

Let  $v = n + x$  giving

The results of the previous sections are all in terms of the L.S.T.'s of various distributions and their first two moments. Inversion of these L.S.T.'s is possible only in the simplest cases. A feasible approach to this problem involves the inversion of power series expansions in  $s$  of the quantities of interest. Kshirsagar and Gupta (1967) take this approach and show how one obtains asymptotic results in terms of the basic quantities  $(m, \bar{a}, \bar{\sigma})$  rather than in terms of recurrence-time distributions as done by Pyke (1961b). The details of the following argument can be found in Kshirsagar and Gupta (1967). Referring to the results of the previous sections, the reader will observe that the quantity  $I^{-1}(s)$  appears often in the L.S.T.'s of the moments and distributions derived. In the

### 7. Asymptotic Results

where  $\beta = 1$  if  $u = k$  or if  $v = k$ .

(1.6.12)

$$c_{I,uv}^0(s) = \sum_{j,k} \left\{ I^{-1} \left( \left( \beta_{jk}^0(s) - \beta_{jk}^0(s) \right) \right) \right\} \left( c_{k,uv}^0(s) + \beta_{m,uv}^0(s) \right),$$

and

$$r(s_0, s) = \frac{s_0}{s-s_0} m(s) - m(s_0) \quad (1.6.11)$$

$$m(s_0, s) = \frac{s_0}{s-s_0} m(s) - m(s_0) \quad (1.6.10)$$

details one obtains

order moments of the  $N_j(t_0, t)$  by differentiating (1.6.9). Omitting Kshirsagar and Gupta (1969b) find the L.S.T.'s of the first and second

(1.6.9)

$$\bar{v}(z, s_0, s) = \frac{s_0}{s-s_0} \bar{e} + \frac{s_0}{s-s_0} \left( I^{-1} \left( \beta(s) - \beta(s_0) \right) \right) \left( \bar{v}(z; s) - \bar{e} \right).$$

or in matrix notation

$$(1.7.3) \quad \frac{s}{m-1} \prod_{l=1}^m (1 + s/\lambda_l) = \frac{|I - P_0 + sP_1|}{1} = \frac{|P_1| \cdot |P_1^{-1}(I - P_0) + sI|}{1}$$

In order to evaluate this expression we first need an expression for  $(I - P_0 + sP_1)^{-1}$ . NOW

$$(1.7.2) \quad (I - P_0 + sP_1)^{-1} \times (I - P_0 + sP_1) = I$$

$$(I - P_0 + sP_1)^{-1} = [I - \frac{1}{2}s^2(I - P_0 + sP_1)^{-1}P_2 - \frac{3}{1}sP_3 + \frac{12}{1}s^2P_4 + o(s^2)]^{-1}$$

that where the matrix  $P_k = (P_{ij}^{(k)})$ . Hence from the above we immediately have

$$(I - P_0 + sP_1)^{-1} = [I - \frac{1}{2}s^2(I - P_0 + sP_1)^{-1}P_2 - \frac{3}{1}sP_3 + \frac{12}{1}s^2P_4 + o(s^2)]^{-1}$$

$$(I - P_0 + sP_1)^{-1} = I - \frac{1}{2}s^2P_2 + \frac{6}{1}s^2P_3 - \frac{24}{1}s^4P_4 + o(s^4)$$

In matrix notation the first few terms provide a power series in s of

$$(1.7.1) \quad q_{ij}^{(k)}(s) = p_{ij}^{(k)}(s) - \frac{1}{2}s^2 p_{ij}^{(k)}(s) + \frac{3i}{3} p_{ij}^{(k)}(s) + \dots$$

Let  $\mu_{ij}^{(k)}$  denote the  $k^{\text{th}}$  order moment of  $F_{ij}^{(k)}(t)$  and write the previous section. Recall that  $Q_{ij}^{(k)}(t) = p_{ij}^{(k)}(t)$  and so  $q_{ij}^{(k)}(s) = p_{ij}^{(k)}(s)$ .

Following argument, due to Kshirsagar and Gupta (1967), we derive an asymptotic expression for  $(I - P_0 + sP_1)^{-1}$  and show how the resulting expression is used to obtain limiting results for several cases discussed in

where  $\lambda_1, \lambda_2, \dots, \lambda_{m-1}, 0$  are the eigenvalues of  $P_1^{-1}(I - P_0)$ . The root zero is of multiplicity one, since  $\sum_{j=1}^m p_{1j} = 1, j = 1, 2, \dots, m$  implies that  $(I - P_0)\bar{e} = 0$ . Also  $1/\alpha = |P_1| \lambda_1 \lambda_2 \dots \lambda_{m-1}$ . The first few terms in (1.7.3) are

$$\left\{ 1 - s \sum_{i=1}^{m-1} \lambda_i + s^2 \sum_{i=1}^{m-1} \sum_{j=1}^{m-1} \lambda_i \lambda_j + \sum_{i < j} \lambda_i \lambda_j \right\}$$

$$- s^3 \sum_{i=1}^{m-1} \sum_{j=1}^{m-1} \sum_{k=1}^{m-1} \lambda_i \lambda_j \lambda_k + \sum_{i < j < k} \lambda_i \lambda_j \lambda_k + o(s^3)$$

$$(1.7.4) \quad \left\{ 1 - a_1 s + a_2 s^2 - a_3 s^3 + o(s^3) \right\}$$

Rewriting  $\text{adj}(I - P_0 + sP_1)$  in powers of  $s$  as

$$(1.7.5) \quad \sum_{r=0}^{m-1} s^r H_r, \quad \text{adj}(I - P_0 + sP_1) = \sum_{r=0}^{m-1} s^r H_r,$$

where  $H_r$  is an  $m \times m$  matrix. Kshirsagar and Gupta (1967) have shown that

$$(1.7.6) \quad H_0 P_1 H_0 = \bar{e} v' P_1 \bar{e} v', \quad k_1 = \bar{v}' P_1 \bar{e}, \quad k_1 = \bar{v}' P_1 \bar{e}$$

where  $v_i$  is the cofactor of any element in the  $i$ th row of  $(I - P_0)$  and

$$(1.7.7) \quad H_0 = \bar{e} \bar{v}' = (\bar{v}', \bar{v}', \dots, \bar{v}')$$

From (1.7.6) and (1.7.7) one easily shows that

$$H_0 P_1 H_0 = k_1 H_0, \quad k_i = \bar{v}' P_1 \bar{e}, \quad i = 2, 3, \dots$$

Combining (1.7.4) and (1.7.5) gives  $(I - P_0 + sP_1)^{-1}$  which when substituted into (1.7.2) gives, after a little algebra,

corresponding cofactor and independent of the initial state.  
 Note that, in the limit, the mean number of visits is proportional to the

$$(1.7.13) \quad \lim_{t \rightarrow \infty} \frac{1}{t} M(t) = \alpha_H^0 \cdot$$

(1.7.12) we have that

chapters, however, we shall make extensive use of these equations. From point no further examples using (1.7.8) or (1.7.10) are given; in later and Gupta (1967) give this result, which is not reproduced here. At this sion to obtain an asymptotic expression for  $E[N_j^i(t) \{N_j^i(t) - 1\}]$ . Kshirsagar (1.7.8) can also be used in (1.4.14) and inverting the resulting expres-

$$(1.7.12) \quad M(t) = [M_{ij}^i(t)] = \alpha_H^0 + \alpha_H^1 + \alpha \left( \frac{2}{1} \alpha k_2 - a_1^1 \right) H^0 - I + o(1) \cdot$$

for large  $t$ ,

so that we may substitute (1.7.8) into the above and invert to obtain,

$$(1.7.11) \quad m(s) = (I - q(s))^{-1} q(s) = (I - q(s))^{-1} (I - I) \cdot$$

Recall from Section 1.4 that

$$(1.7.10) \quad \left\{ (I - q(s))^{-1} \right\}_{ij} = \frac{1}{s} v_j + a v_j + \alpha h_{ij} + s \delta_{ij} + o(s) \cdot$$

$(I - q(s))^{-1}$  is given by

where  $a = \alpha \left( \frac{2}{1} \alpha k_2 - a_1^1 \right)$  and  $L$  is an  $m \times m$  matrix. The  $(i, j)$ th element of

$$(1.7.9) \quad \frac{s}{\alpha} H^0 + \alpha_H^1 + \alpha_H^0 + sL + o(s) \cdot$$

$$(1.7.8) \quad - \frac{1}{2} \alpha^2 (a_1^1 k_2 + \frac{3}{1} k_3) H^0 + o(s)$$

$$+ \alpha s \{ H^2 - a_1^1 H^1 + a_2^2 H^0 + \frac{1}{2} \alpha (H^0 P^2 H^1 + H^1 P^2 H^0) \}$$

$$(I - q(s))^{-1} = \frac{s}{\alpha} H^0 + \alpha_H^1 + \alpha \left( \frac{2}{1} \alpha k_2 - a_1^1 \right) H^0$$

$N_{i \cdot k}^k(t)$  and  $N_{\cdot k}^k(t)$  being the row and column sums of  $N(t)$  and where  $i$  and  $j$

$$(2.1.2) \quad N_{i \cdot k}^k(t) - N_{\cdot k}^k(t) = \delta_{ik} - \delta_{jk} \quad , \quad k = 1, 2, \dots, m \quad ,$$

is constrained by

state  $i$  except for the initial and final states. Hence the matrix  $N(t)$

Observe that every entry into state  $i$  must be followed by an exit from

$$(2.1.1) \quad = \Pr\{N_i^i(t) = k \quad , \quad \bar{J}(t) = \bar{e}_j \quad \} .$$

$$W_{ij}^{ij}(k;t) = \Pr\{N(t) = k \quad , \quad \bar{J}(t) = \bar{e}_j \mid \bar{J}(0) = \bar{e}_i \quad \}$$

integers and define

of transition counts. Further, let  $k$  be an  $m \times m$  matrix of non-negative

Let  $N(t)$  be the  $m \times m$  matrix of the  $N_{ij}^{ij}(t)$ . We shall call  $N(t)$  the matrix

state  $j$  in the interval  $(0, t]$  for the M.R.P. specified by  $(m, \bar{a}, \bar{Q})$  and

Let  $N_{ij}^{ij}(t)$  be the number of one-step transitions from state  $i$  to

### 1. Derivation

a generalization of both Whittle's result and Pyke's results.

tinuous and derive the analogous distribution for M.R.P.'s. Thus we provide

state  $j$  to state  $k$ , for finite Markov chains. Here we consider time con-

$[n^{jk}]$  whose  $(j, k)$ th element is the number of one-step transitions from

P. Whittle (1955) derived the probability distribution of the matrix

## TRANSITION FREQUENCIES IN M.R.P.'S

### CHAPTER II



are the initial and final states respectively. For a matrix  $k$  satisfying

(2.1.2) we have that

$$W_{ij}^{1j}(k,t) = \sum_{m=1}^m \delta_{im}^{1j}(t) * W_{mj}^{1j}(k - e_{im}^{1j}; t) ,$$

(2.1.3)

$$W_{ij}^{1j}(0,t) = \delta_{ij}^{1j}(1 - H_{ij}^{1j}(t)) .$$

Define

$$W_{ij}^{1j}(z,s) = \int_0^{\infty} \sum_{k \geq 0} e^{-st} \delta_{ij}^{1j}(k,t) W_{mj}^{1j}(k,t) \Pi_{z_{ij}^{1j}}^{k_{ij}^{1j}} , \quad (2.1.4)$$

where  $z = (z_{ij}^{1j}) , |z_{ij}^{1j}| \leq 1 , i, j = 1, 2, \dots, m ; i.e., W_{ij}^{1j}(z,s)$  is the L.S.T. of the p.g.f. of  $W_{ij}^{1j}(k,t)$ . Using (2.1.4) in (2.1.3) gives

$$W_{ij}^{1j}(z,s) = \delta_{ij}^{1j}(1 - h_{ij}^{1j}(s)) + \sum_{m=1}^m \delta_{im}^{1j}(s) z_{mj}^{1j} W_{mj}^{1j}(z,s) ,$$

which in matrix notation becomes

$$W(z,s) = I - h(s) + (q(z) \square W(z,s)) ,$$

where  $A \square B$  is a matrix with  $(i,j)$ th element  $a_{ij} b_{ij}$ . Solving for  $W(z,s)$  gives

$$W(z,s) = (I - h(s))^{-1} (q(z) \square W(z,s)) . \quad (2.1.5)$$

Summing the rows of  $W(z,s)$  gives

$$\bar{w}(z,s) = (I - q(s))^{-1} (I - q(s)) \bar{e} \quad (2.1.6)$$

which is the L.S.T. of the p.g.f. of

$$W_{ij}^{1j}(k,t) = Pr\{N_{ij}^1(t) = k\} \quad (2.1.7)$$

if  $N(t)$  satisfies (2.1.2) and zero otherwise. The analogous result for  $w^T(z;s)$  is obtained by summing (2.1.10) with respect to the state  $j$ . Hence the only contribution from summing (2.1.10) with respect to the state  $j$  arises when  $j = k$  where  $k$  is the solution of the system of equations (2.1.2) is unique given the matrix  $N(t)$  and the initial state  $i$ . Martin (1967, pg. 119) gives a simple proof showing that the solution for

$$(2.1.10) \quad \prod_{N, k, i} \frac{\prod_{N, k, \lambda} \prod_{N, k, \lambda} (s)}{\prod_{N, k, \lambda} (s) \prod_{N, k, \lambda} (s)} \left( \prod_{N, k, \lambda} (s) \right) \left( 1 - h_k(s) \right)$$

With the aid of the lemma the coefficient of  $\prod_{N, k, \lambda} z_{k, \lambda}$  in  $w^T(z;s)$  is

$$(2.1.9) \quad \left\{ \begin{array}{l} \delta_{i, j} \\ \delta_{i, j} - n_{i, j} / n_{i, j} \end{array} \right. = n_{i, j}^* \quad ; \quad n_{i, j} > 0$$

of the matrix  $N^* = (n_{i, j}^*)$ , defined by and is zero otherwise. Here  $n_{i, j}^*$  is the cofactor of the  $(j, i)$ th element

$$n_{i, j} - n_{i, k} = \delta_{i, k} - \delta_{j, k} \quad , \quad k = 1, 2, \dots, m \quad ,$$

if  $n_{k, \lambda}$  are non-negative integers such that

$$(2.1.8) \quad \prod_{N, k, i} \frac{\prod_{N, k, \lambda} n_{k, \lambda}}{\prod_{N, k, i} n_{k, i}} \quad ,$$

in the  $(i, j)$ th element of the  $m \times m$  matrix  $(I - A)^{-1}$ , where  $A = (a_{i, j})$ , is

$$\prod_{N, k, \lambda} a_{k, \lambda}$$

Lemma. [Whittle (1955)]. The coefficient of

(2.2.2)

$$= \sum_{k=1}^m \delta_{\alpha}^{1k} (s) \delta_{\beta}^{1k} + \sum_{k=1}^m \delta_{\alpha}^{1k} (s) \delta_{\beta}^{1k} (s) \delta_{\alpha}^{1k} (s)$$

$$= \sum_{k=1}^m \delta_{\alpha}^{1k} (s) \delta_{\beta}^{1k} + \sum_{k=1}^m \delta_{\alpha}^{1k} (s) \delta_{\beta}^{1k} (s) \delta_{\alpha}^{1k} (s)$$

Taking L.S.T.'s of the expectation on both sides,

(2.2.1)

$$\left. \begin{aligned} & 0, \text{ if there is no transition in } (0, t]. \\ & \delta_{\alpha}^{1k} (t) + N_k^{\alpha} (t) - x_k^1, \quad \delta_{\beta}^{1k} = \bar{e}_k \end{aligned} \right\} = N_k^{\alpha} (t)$$

more direct approach. For the first method note that the second method involves differentiation of (2.1.6) and its perhaps the may also be found in a recent paper by Kshirsagar and Wysocki (1969a). of a method developed by Martin (1967, pg. 122ff) for Markov Chains and give two independent derivations of the moments. The first is an extension the L.S.T.'s of these moments by  $\bar{m}^{jk}(s)$  and  $\bar{c}^{jk, \lambda n}(s)$  respectively. We and let  $M^{jk}(t)$  and  $C^{jk, \lambda n}(t)$  be the  $m \times 1$  vectors of these moments. Denote

$$\left\{ \begin{aligned} M^{jk}(t) &= E \left\{ N^{jk}(t) \right\} \\ C^{jk, \lambda n}(t) &= E \left\{ N^{jk}(t) N^{\lambda n}(t) \right\} \end{aligned} \right.$$

Define

2. Moments of  $N^1(t)$

where  $k$  is the unique solution of (2.1.2).

(2.1.11)

$$N_k^{\alpha} = \frac{\alpha}{\alpha - \beta} \frac{\beta}{\beta - \alpha} \left( \frac{\beta}{\beta - \alpha} \right)^{\alpha} (s) \left( \frac{\alpha}{\alpha - \beta} \right)^{\beta} (s) \left( 1 - h_k(s) \right)$$

(2.1.2). Hence the coefficient of  $z_k^{\lambda}$  in  $w^1(z; s)$  is

or in matrix notation

$$(2.2.3) \quad \bar{m}^{\alpha\beta}(s) = q^{\alpha\beta}(s)e^{-\alpha} + q(s)\bar{m}^{\alpha\beta}(s) \cdot$$

Solving for  $\bar{m}^{\alpha\beta}(s)$  gives

$$(2.2.4) \quad \bar{m}^{\alpha\beta}(s) = q^{\alpha\beta}(s) \left( I - q(s) \right)^{-1} e^{-\alpha}$$

whose  $i$ th element is

$$(2.2.5) \quad \bar{m}_i^{\alpha\beta}(s) = q^{\alpha\beta}(s) \left( I - q(s) \right)^{-1} \left\{ \begin{matrix} \alpha \\ i \end{matrix} \right\} \cdot$$

Just as in (2.2.1) we can write

$$N_i^{\alpha\beta}(t) N_i^{\gamma\delta}(t) = \delta^{\alpha\gamma} \delta^{\beta\delta} \delta^{\alpha\beta} \delta^{\gamma\delta} + \delta^{\alpha\gamma} \delta^{\beta\delta} \delta^{\alpha\delta} N_i^{\alpha\beta}(t - x_1) + \delta^{\alpha\gamma} \delta^{\beta\delta} \delta^{\gamma\delta} N_i^{\alpha\beta}(t - x_1)$$

$$(2.2.6) \quad + \delta^{\alpha\gamma} \delta^{\beta\delta} \delta^{\alpha\delta} N_i^{\alpha\beta}(t - x_1) + N_i^{\alpha\beta}(t - x_1) N_i^{\gamma\delta}(t - x_1)$$

if  $\bar{J}_i = e_k$  for  $k = 1, 2, \dots, m$  and is zero if there is no transition

in  $(0, t]$ . Taking L.S.T.'s of the expectations on both sides of (2.2.6)

one can write in matrix notation that

$$(2.2.7) \quad \bar{c}^{\alpha\beta, \gamma\delta}(s) = \delta^{\alpha\gamma} \delta^{\beta\delta} \bar{m}^{\alpha\beta}(s) + \bar{m}^{\alpha\beta}(s) \bar{m}^{\gamma\delta}(s) + \bar{m}^{\alpha\beta}(s) \bar{m}^{\gamma\delta}(s) \cdot$$

whose  $i$ th element is

$$(2.2.8) \quad \bar{c}_i^{\alpha\beta, \gamma\delta}(s) = \delta^{\alpha\gamma} \delta^{\beta\delta} \bar{m}_i^{\alpha\beta}(s) + \bar{m}_i^{\alpha\beta}(s) \bar{m}_i^{\gamma\delta}(s) + \bar{m}_i^{\alpha\beta}(s) \bar{m}_i^{\gamma\delta}(s) \cdot$$

The second method follows from the appropriate differentiation of

$$(2.2.9) \quad \bar{w}(z; s) = (I - q(s) \square z) \left( I - q(s) \right)^{-1} \bar{e}$$

and evaluation at  $z = e$ . Rather than proceeding from (2.2.9) directly,

it is more convenient and algebraically simpler to find an expression for the L.S.T. of the marginal distribution of  $(N_i^{\alpha\beta}(t), N_i^{\gamma\delta}(t))$  and then find

$$(2.2.15) \quad \bar{c}_{\alpha\beta, \gamma\delta}^i(s) = -\bar{m}_i^{\gamma\delta}(s) \bar{m}_i^{\alpha\beta}(s) + \bar{m}_i^{\alpha\beta}(s) \bar{m}_i^{\gamma\delta}(s) \quad .$$

whose  $i$ th element is

$$(2.2.14) \quad = \bar{q}^{\alpha\beta}(s) \bar{q}^{\gamma\delta}(s) - \bar{q}^{\gamma\delta}(s) \bar{q}^{\alpha\beta}(s) + \bar{q}^{\gamma\delta}(s) \bar{q}^{\alpha\beta}(s) - \bar{q}^{\alpha\beta}(s) \bar{q}^{\gamma\delta}(s) \quad ,$$

$$\bar{c}_{\alpha\beta, \gamma\delta}^i(s) = \left. \frac{\partial^2 \bar{w}(z^*, s)}{\partial z_1^{\alpha\beta} \partial z_2^{\gamma\delta}} \right|_{z_1=z_2=1}$$

product moments  $\bar{c}_{\alpha\beta, \gamma\delta}^i(t)$  has L.S.T. given by

of  $E\{N_{\beta}^i(t)\}$  derived by Pyke (1961b). For  $\alpha \neq \gamma$  or  $\beta \neq \delta$  the vector of which is the L.S.T. of the matrix of means with  $(i, \beta)$ th element the L.S.T.

$$(2.2.13) \quad \left( \bar{m}_i^{\alpha\beta}(s) \right) = \left( \sum_{\alpha=1}^m \bar{m}_i^{\alpha\beta}(s) \right) \left( \bar{q}^{\alpha\beta}(s) - \bar{q}^{\beta\alpha}(s) \right) \quad ,$$

in agreement with (2.2.4). Note that

$$(2.2.12) \quad = \bar{q}^{\alpha\beta}(s) \bar{q}^{\beta\alpha}(s) - \bar{q}^{\beta\alpha}(s) \bar{q}^{\alpha\beta}(s) \quad ,$$

$$\bar{m}_i^{\alpha\beta}(s) = \left. \frac{\partial \bar{w}(z^*, s)}{\partial z_1^{\alpha\beta}} \right|_{z_1=z_2=1} = \bar{q}^{\alpha\beta}(s) \bar{q}^{\beta\alpha}(s) - \bar{q}^{\beta\alpha}(s) \bar{q}^{\alpha\beta}(s) \quad ,$$

Then the L.S.T. of the vector of conditional means  $\bar{m}_i^{\alpha\beta}(s)$  is given by

$$(2.2.11) \quad \bar{w}(z^*, s) = \left[ \bar{q}^{\alpha\beta}(s) - \bar{q}^{\beta\alpha}(s) \right] \bar{q}^{\alpha\beta}(s) e^{\gamma\delta} - \left[ \bar{q}^{\alpha\beta}(s) - \bar{q}^{\beta\alpha}(s) \right] \bar{q}^{\gamma\delta}(s) e^{\gamma\delta} \quad ,$$

substituting (2.2.10) into (2.2.9) one easily obtains

$$(2.2.10) \quad Z = Z^* = E + (z_1 - 1)e^{\alpha\beta} + (z_2 - 1)e^{\gamma\delta} \quad .$$

moments from the resulting expression. To that end set

For  $\alpha = \gamma$  and  $\beta = \delta$  we obtain the L.S.T. of  $E\{N_1^{\alpha\beta}(t) - 1\}$  which is  $2m_1^{\alpha\beta} m_2^{\alpha\beta}$  so that the L.S.T. of  $E\{N_1^{\alpha\beta}(t)\}$  is

$$(2.2.16) \quad c_1^{\alpha\beta, \alpha\beta} = m_1^{\alpha\beta} (zm_2^{\alpha\beta} + 1) .$$

Combining (2.2.15) and (2.2.16) we have that for any  $i, \alpha, \beta, \gamma, \delta$  that

$$(2.2.17) \quad c_1^{\alpha\beta, \gamma\delta}(s) = \delta^{\alpha\gamma} \delta^{\beta\delta} m_1^{\alpha\beta}(s) + m_1^{\alpha\beta}(s) m_1^{\gamma\delta}(s) + m_1^{\beta\delta}(s) m_1^{\alpha\gamma}(s) ,$$

in agreement with (2.2.8).

### 3. Cumulative Processes

In this chapter we have associated the holding time  $X_n$  with the

transition from state  $J^{n-1}$  to state  $J^n$ . We can generalize this concept

by considering, in place of  $X_n$ , the real-valued function  $g(J^{n-1}, J^n, X_n)$ .

For  $n = 1, 2, \dots$  we assume that the  $g$ -functions are independent and

identically distributed. Then the random variable defined at time  $t$  by

$$(2.3.1) \quad W_g(t) = \begin{cases} 0 & , N(t) = 0 \\ \sum_{n=1}^{N(t)} g(J^{n-1}, J^n, X_n) & , N(t) > 0 \end{cases}$$

where  $N(t) = \sum_{j=1}^J N_j(t)$ , gives rise to the stochastic process  $\{W_g(t) : t \geq 0\}$  which is known as a cumulative process. The reader is referred to Cox

(1962, Chapter 8) for examples of how cumulative processes arise in appli-

cations.

Let

$$(2.3.2) \quad \xi_{T_j}(p) = \int_{-\infty}^0 e^{-p g(t, j, x)} dO_{T_j}^x(x)$$

be the moment generating function of  $g(t, j, x)$ . Then, since the  $g$ -functions are independent, we can write

$$(2.3.3) \quad \mathbb{E} \left\{ e^{-pW^g(t)} \mid N^j(t) \right\} = \mathbb{E} \left\{ \prod_{j,k} \left[ \xi^{jk}(p) \right]^{N^j_k(t)} \right\} .$$

So that

$$(2.3.4) \quad \mathbb{E} \left\{ e^{-pW^g(t)} \mid \bar{J}(0) = \bar{e}_1 \right\} = \mathbb{E} \left\{ \prod_{j,k} \left[ \xi^{jk}(p) \right]^{N^j_k(t)} \right\} .$$

Comparing (2.3.4) with (2.1.4) we see that this is nothing but the p.g.f. of  $N^j_k(t)$  with argument  $\xi^{jk}(p)$  hence the l.s.f. of  $W^g(t) \mid \bar{J}(0) = \bar{e}_1$  is the  $i$ th element of the vector

$$(2.3.5) \quad [I - q(s) \square \xi(p)]^{-1} (I - p(s)) \bar{e} ,$$

where  $\xi(p) = \left( \xi^{jk}(p) \right)$ . To find the moments of  $W^g(t)$  one may proceed by

differentiating (2.3.5), but a more direct approach is available. To find

the mean note that

$$\mathbb{E} \left\{ W^g(t) \mid \bar{J}(0) = \bar{e}_1 \right\} = \mathbb{E} \left\{ W^g(t) \mid N(t) , \bar{J}(0) = \bar{e}_1 \right\}$$

$$(2.3.6) \quad = \mathbb{E} \left\{ \sum_{j,k} N^j_k(t) \int_0^\infty g(j,k,x) dF^{jk}(x) \right\} .$$

Following the notation in Pyke and Moore (1968) we set

$$(2.3.7) \quad A^{jk} = \int_0^\infty g(j,k,x) dQ^{jk}(x) , \quad A_j = \sum_k A^{jk} ,$$

$$(2.3.8) \quad B^{jk} = \int_0^\infty g_2(j,k,x) dQ^{jk}(x) , \quad B_j = \sum_k B^{jk} .$$

Then

$$\mathbb{E} \left\{ W^g(t) \mid \bar{J}(0) = \bar{e}_1 \right\} = \mathbb{E} \left\{ \sum_{j,k} \frac{A^{jk}}{P^{jk}} N^j_k(t) \right\} = \sum_{j,k} \frac{A^{jk}}{P^{jk}} M^j_k(t) .$$

(2.3.9)

$$E\{N_{jk}^T(t)\} = P_{jk} \{ \alpha v_j t + a v_j + \alpha(h_{1j} - v_j \mu_{jk}) + o(1) \} \quad (2.4.4)$$

whose inverse transform is, for large  $t$ ,

$$m_{jk}^T(s) = P_{jk} \left\{ \frac{\alpha}{s} v_j + a v_j + \alpha(h_{1j} - v_j \mu_{jk}) + s \left( \lambda_{1j} - \mu_{jk} \{ a v_j + \alpha h_{1j} \} + \frac{1}{2} \alpha v_j \mu_{jk} \right) + o(s) \right\} \quad (2.4.3)$$

Using (2.4.1) and (2.4.2) in (2.2.5) one obtains

$$q_{jk}^T(s) = P_{jk} \left\{ 1 - s \mu_{jk} + \frac{1}{2} s^2 \mu_{jk}^2 + o(s^2) \right\} \quad (2.4.2)$$

and  $q_{jk}^T(s)$  could be written in the form

$$\left\{ (I - q(s))^{-1} \right\}_{1j} = \frac{s}{\alpha} v_j + a v_j + \alpha h_{1j} + s \lambda_{1j} + o(s) \quad (2.4.1)$$

of transition counts. Recall that in Section 1.7 we had shown that values of the first few moments of the distribution of  $N(t)$ , the matrix The results of Section 1.7 can be used to determine the asymptotic

#### 4. Asymptotic Results

$$= \sum_{j,k} \left[ \frac{B_{jk}}{P_{jk}} - \frac{P_{jk}^2}{A_{jk}^2} \right] M_{jk}^T(t) + \sum_{j,k} \lambda_{j,n} \left[ \frac{A_{jk}^2}{P_{jk}^2} \text{Cov} \left\{ N_{jk}^T(t), N_{jn}^T(t) \right\} \right] \quad (2.3.10)$$

$$\text{Var} \left\{ W_g^T(t) | \bar{J}(0) = \bar{e}_{-1} \right\} = E \left\{ \sum_{j,k} \left[ \frac{B_{jk}}{P_{jk}} - \frac{P_{jk}^2}{A_{jk}^2} \right] N_{jk}^T(t) \right\} + \text{Var} \left\{ \sum_{j,k} \frac{A_{jk}^2}{P_{jk}^2} N_{jk}^T(t) \right\}$$

we obtain

$$\text{Var} \left\{ W_g^T(t) | \bar{J}(0) = \bar{e}_{-1} \right\} = E \text{Var} \left\{ W_g^T(t) | N^T(t) \right\} + \text{Var} E \left\{ W_g^T(t) | N^T(t) \right\}$$

From the identity



$$\lim_{t \rightarrow \infty} \frac{1}{t} \text{Var} \{ N_{jk}^i(t) \} = \alpha v_{jp}^{jk} + 2\alpha v_{jp}^{jk} (av_j - \alpha v_{jp}^{jk} + \alpha h_{kj}) \quad (2.4.10)$$

and

$$\lim_{t \rightarrow \infty} \frac{1}{t} \text{Cov} \{ N_{jk}^i(t), N_{ln}^i(t) \} = P_{jk}^{ln} [2\alpha v_{jp}^{jk} - \alpha^2 v_{jp}^{jk} (v_{h_{kl}} + v_{h_{nj}})] + \alpha^2 (v_{h_{kl}} + v_{h_{nj}}) \quad (2.4.9)$$

Taking the appropriate limits in (2.4.7) and (2.4.8) gives

$$\text{Var} \{ N_{jk}^i(t) \} = \alpha v_{jp}^{jk} t + \alpha v_{jp}^{jk} t (2\alpha v_j - 2\alpha v_{jp}^{jk} + 2\alpha h_{kj}) + o(t) \quad (2.4.8)$$

and

$$\text{Cov} \{ N_{jk}^i(t), N_{ln}^i(t) \} = \delta_{jl} \delta_{kn} \alpha v_{jp}^{jk} t + P_{jk}^{ln} [2\alpha v_{jp}^{jk} - \alpha^2 v_{jp}^{jk} (v_{h_{kl}} + v_{h_{nj}})] + o(t) \quad (2.4.7)$$

From the inverses of (2.4.6) and (2.4.3) we get, for large  $t$ ,

$$\begin{aligned} \frac{1}{t} \text{Cov} \{ N_{jk,ln}^i(s) \} &= \frac{1}{2} \alpha^2 v_{jp}^{jk} v_{lp}^{ln} + \frac{1}{s} \delta_{jl} \delta_{kn} P_{jk}^{ln} \\ &+ \frac{1}{s} P_{jk}^{ln} [4\alpha v_{jp}^{jk} + \alpha^2 (v_{h_{kl}} + v_{h_{lj}})] \\ &- 2\alpha^2 v_{jp}^{jk} (v_{h_{kl}} + v_{h_{nj}}) + \alpha^2 (v_{h_{kl}} + v_{h_{nj}}) + o(1/s) \end{aligned} \quad (2.4.6)$$

expect. Using (2.4.3) in (2.2.17) gives

which is independent of the initial state of the system as one would

$$\lim_{t \rightarrow \infty} \frac{1}{t} E \{ N_{jk}^i(t) \} = \alpha v_{jp}^{jk} \quad (2.4.5)$$

Hence

The reader will note that the above expressions are independent of the

initial state of the system. From (2.4.4) we also have

$$\lim_{t \rightarrow \infty} \frac{t}{1} E\{W^g(t) | \bar{J}(0) = \bar{e}_1\} = \sum_{j,k} \frac{A_{jk}^j}{P_{jk}^j} \alpha P_{jk}^j \alpha V_{jk}^j = \sum_{j,k} \alpha V_{jk}^j A_{jk}^j \quad (2.4.11)$$

The asymptotic variance is formed as follows:

$$\lim_{t \rightarrow \infty} \frac{t}{1} \text{Var}\{W^g(t) | \bar{J}(0) = \bar{e}_1\} = \bar{e}_1 \left\{ \sum_{j,k} \alpha V_{jk}^j A_{jk}^j - \left[ \sum_{j,k} \frac{A_{jk}^j \alpha V_{jk}^j}{P_{jk}^j} \right]^2 + \sum_{j,k} \alpha V_{jk}^j A_{jk}^j \left[ \sum_{n} \delta_{jk} \delta_{kn} \alpha V_{jn}^j P_{jk}^j \right] \right\}$$

$$= \alpha \sum_{j,k} V_{jk}^j B_{jk}^j - \alpha \sum_{j,k} \frac{P_{jk}^j}{V_{jk}^j} A_{jk}^j + \alpha \sum_{j,k} \frac{P_{jk}^j}{V_{jk}^j} A_{jk}^j + \alpha \sum_{j,k} \sum_{n} A_{jk}^j \alpha V_{jn}^j \{2\alpha V_{jk}^j + P_{jk}^j P_{jn}^j [2\alpha V_{jk}^j - \alpha^2 V_{jk}^j V_{jn}^j + \mu_{jn}^j] + V_{jn}^j \alpha V_{jk}^j\}$$

$$= \alpha \sum_{j,k} V_{jk}^j B_{jk}^j - \alpha \sum_{j,k} \frac{P_{jk}^j}{V_{jk}^j} A_{jk}^j + \alpha \sum_{j,k} \frac{P_{jk}^j}{V_{jk}^j} A_{jk}^j + 2\alpha \sum_{j,k} A_{jk}^j \alpha V_{jn}^j \{2\alpha V_{jk}^j + P_{jk}^j P_{jn}^j [2\alpha V_{jk}^j - \alpha^2 V_{jk}^j V_{jn}^j + \mu_{jn}^j] + V_{jn}^j \alpha V_{jk}^j\}$$

$$= \alpha \sum_{j,k} V_{jk}^j B_{jk}^j + \alpha \left[ \sum_{j,k} \frac{P_{jk}^j}{V_{jk}^j} A_{jk}^j - \sum_{j,k} \frac{P_{jk}^j}{V_{jk}^j} A_{jk}^j \right] + 2\alpha \sum_{j,k} A_{jk}^j \alpha V_{jn}^j$$

$$= \alpha \sum_{j,k} V_{jk}^j B_{jk}^j + \alpha \left[ \sum_{j,k} \frac{P_{jk}^j}{V_{jk}^j} A_{jk}^j - \sum_{j,k} \frac{P_{jk}^j}{V_{jk}^j} A_{jk}^j \right] + 2\alpha \sum_{j,k} A_{jk}^j \alpha V_{jn}^j$$

$$- 2\alpha^2 \sum_{j,k} \sum_{\lambda} A_{jk}^j \alpha V_{jk}^j (V_{jk}^j \alpha V_{jk}^j - h_{jk}^{\lambda}) \quad (2.4.12)$$

(3.1.6) 
$$W_I^T(k; t) = \text{Pr}\{N_I^T(t) = k\}$$

Similar notation will be used for the I.S.T.'s of  $R_I^{jk}(t_0, t)$  and  $C_I^{jk, \lambda n}(t_0, t)$ . Letting

(3.1.5) 
$$m_I^{jk}(s_0, s) = \int_{-\infty}^0 e^{-s_0 t_0} m_I^{jk}(t_0, t) dt_0 \cdot$$

(3.1.4) 
$$M_I^{jk}(s_0, s) = \int_{-\infty}^0 e^{-s t_0} M_I^{jk}(t_0, t) dt_0 \cdot$$

The following I.S.T.'s will also be used:

(3.1.3) 
$$C_I^{jk, \lambda n}(t_0, t) = E\{N_I^{jk}(t_0, t) N_I^{\lambda n}(t_0, t)\} \cdot$$

and

(3.1.2) 
$$R_I^{jk}(t_0, t) = E\{N_I^{jk}(t_0, t) [N_I^{jk}(t_0, t) - 1]\} \cdot$$

(3.1.1) 
$$M_I^{jk}(t_0, t) = E\{N_I^{jk}(t_0, t)\} \cdot$$

$(t_0, t_0 + t]$ . Further we define

the number of direct transitions from state  $i$  to state  $j$  in the interval  $(t_0, t_0 + t]$  is derived. Specifically, let  $N_I^{jk}(t_0, t)$  denote the distribution of the matrix of transition counts in the

In this chapter the results of Section 1.6 and Chapter II are com-

1. Derivation

FOR ARBITRARY INTERVALS

DISTRIBUTION OF TRANSITION FREQUENCIES

CHAPTER III

(3.1.11)

$$+ \frac{s_0}{1} - \int_{-1}^1 \left\{ \left( I - \left( \begin{smallmatrix} 0 & p \\ s_0 & 0 \end{smallmatrix} \right) \right)^{j_1} \right\} \left( \begin{smallmatrix} 0 & p \\ s_0 & 0 \end{smallmatrix} \right)^{j_2} \left( \begin{smallmatrix} 0 & p \\ s_0 & 0 \end{smallmatrix} \right)^{j_3} \dots$$

$$= \int_{-1}^1 \left\{ \left( I - \left( \begin{smallmatrix} 0 & p \\ s_0 & 0 \end{smallmatrix} \right) \right)^{j_1} \right\} \left( \begin{smallmatrix} 0 & p \\ s_0 & 0 \end{smallmatrix} \right)^{j_2} \left( \begin{smallmatrix} 0 & p \\ s_0 & 0 \end{smallmatrix} \right)^{j_3} \dots$$

$$\int_{-\infty}^0 e^{-s_0 t} w^T(z; t_0; s) dt_0$$

whose L.S.T. with respect to  $t_0$ , using (3.1.9), is

$$(3.1.10) \quad + 1 - \int_{-1}^1 \left\{ \left( I - \left( \begin{smallmatrix} 0 & p \\ s_0 & 0 \end{smallmatrix} \right) \right)^{j_1} \right\} \left( \begin{smallmatrix} 0 & p \\ s_0 & 0 \end{smallmatrix} \right)^{j_2} \dots = e^{-1}$$

$$w^T(z; t_0; s) = \int_{-1}^1 \left\{ \left( I - \left( \begin{smallmatrix} 0 & p \\ s_0 & 0 \end{smallmatrix} \right) \right)^{j_1} \right\} \left( \begin{smallmatrix} 0 & p \\ s_0 & 0 \end{smallmatrix} \right)^{j_2} \dots = e^{-1}$$

write

$N(t_0; t)$  conditional on the initial state. Then, using (3.1.7) we can

Define  $\bar{w}(z; t_0; s)$  as the L.S.T. with respect to  $t$  of the p.g.f. of

$$(3.1.9) \quad \left( \begin{smallmatrix} 0 & p \\ s_0 & 0 \end{smallmatrix} \right)^{j_1} \left\{ \left( I - \left( \begin{smallmatrix} 0 & p \\ s_0 & 0 \end{smallmatrix} \right) \right)^{j_2} \right\} \left( \begin{smallmatrix} 0 & p \\ s_0 & 0 \end{smallmatrix} \right)^{j_3} \dots = e^{-1}$$

and that its L.S.T. is given by

$$(3.1.8) \quad + \int_{t_0}^0 \frac{dx}{p} \delta^{j_1}(t_0; n) \delta^{j_2}(n + x_0)$$

$$\frac{dx}{p} \delta^{j_1}(t_0; n) \delta^{j_2}(n + x_0) = e^{-1} = \left( \begin{smallmatrix} 0 & p \\ s_0 & 0 \end{smallmatrix} \right)^{j_1} \left\{ \left( I - \left( \begin{smallmatrix} 0 & p \\ s_0 & 0 \end{smallmatrix} \right) \right)^{j_2} \right\} \dots$$

Recall from Section 1.6 that the p.d.f. of the first transition after  $t_0$  is

$$(3.1.7) \quad \bar{w}(z; s) = \left( I - \left( \begin{smallmatrix} 0 & p \\ s_0 & 0 \end{smallmatrix} \right) \right)^{-1} \left( \begin{smallmatrix} 0 & p \\ s_0 & 0 \end{smallmatrix} \right)$$

$\bar{w}(k; t)$  is

it has been previously shown that the L.S.T. of the p.g.f. of the vector

$$(3.2.3) \quad \begin{aligned} & \left( \frac{s^0}{1} \bar{m}^{jk} - (s) \bar{m}^{jk} \right) = \\ & \left[ (s) \bar{m}^{jk} + e^{jk} \bar{e}^{jk} - (s) \bar{e}^{jk} \right] \left( (s) \bar{d} - I \right) - \\ & \left( (s) \bar{d} - I \right) \left( (s) \bar{d} - I \right) + \\ & \left( (s) \bar{m}^{jk} - (s) \bar{e}^{jk} \right) \left( (s) \bar{d} - I \right) \left[ \frac{s^0}{1} \bar{m}^{jk} - (s) \bar{m}^{jk} \right] = \\ & \left( (s) \bar{m}^{jk} \right) \left( (s) \bar{d} - I \right) \frac{s^0}{1} + \\ & \left( (s) \bar{e}^{jk} \right) \left( (s) \bar{d} - I \right) \frac{s^0}{1} = (s) \bar{m}^{jk}, \end{aligned}$$

At  $Z = E$  we obtain the L.S.T. of the vector of conditional means

$$(3.2.2) \quad \frac{\partial \bar{w}(Z; s)}{\partial z^{jk}} = \frac{\partial z^{jk}}{(s) \bar{m}^{jk}} + \left( (s) \bar{d} - I \right) \left( (s) \bar{d} - I \right) \frac{s^0}{1} + \left( (s) \bar{d} - I \right) \frac{\partial \bar{w}(Z; s)}{\partial z^{jk}}$$

into (3.1.12) and proceeding as in Section 2.2. However, it is less cumbersome to proceed directly from (3.1.12). To that end

$$(3.2.1) \quad Z^* = Z + E + (z_1 - 1)e^{jk} + (z_2 - 1)e^{kn}$$

One method involves substituting by appropriate differentiation and evaluating these expressions at  $z = e$ . From (3.1.12) we can obtain the L.S.T.'s of the moments of  $N(t, t)$

## 2. Moments

$$(3.1.12) \quad \bar{w}(Z; s) = \frac{s^0}{1} \bar{e} + \left[ \frac{s^0}{1} \bar{e} - (s) \bar{d} - I \right] \left[ \left( (s) \bar{d} - I \right) \bar{w}(Z; s) \right] + (s) \bar{e}^{jk}$$

In matrix notation the vector  $\bar{w}(Z; s)$  is

(3.2.8)

$$c_{jk, \lambda n}^{(s_0, s)} = \delta_{j\lambda} \delta_{kn} m_{jk}^{(s_0, s)} + m_n^{(s_0, s)} + m_k^{(s_0, s)} m_{\lambda n}^{(s_0, s)} + m_{\lambda n}^{(s_0, s)} m_{jk}^{(s_0, s)}$$

In general, for any  $i, j, k, \lambda$ , and  $n$

(3.2.7)

$$r_{jk}^{(s_0, s)} = 2m_{jk}^{(s_0, s)}$$

For  $j \neq \lambda$  or  $k \neq n$ . If  $j = \lambda$  and  $k = n$ , then

(3.2.6)

$$m_n^{(s_0, s)} = m_n^{(s_0, s)} + m_k^{(s_0, s)} m_{\lambda n}^{(s_0, s)}$$

(3.2.5)

$$c_{jk, \lambda n}^{(s_0, s)} = \frac{s_0}{1-s} m_n^{(s_0, s)} [m_{\lambda n}^{(s_0, s)} - \frac{s_0}{1-s} m_{\lambda n}^{(s_0, s)}] + \frac{s_0}{1-s} m_k^{(s_0, s)} m_{\lambda n}^{(s_0, s)} [m_{jk}^{(s_0, s)} - m_{jk}^{(s_0, s)}]$$

whose  $i$ th element is, after some simplification,

(3.2.4)

$$c_{jk, \lambda n}^{(s_0, s)} = \frac{s_0}{1-s} m_k^{(s_0, s)} m_{\lambda n}^{(s_0, s)} \left( \frac{s_0}{1-s} m_{\lambda n}^{(s_0, s)} - m_{\lambda n}^{(s_0, s)} \right) + \frac{s_0}{1-s} m_n^{(s_0, s)} \left( \frac{s_0}{1-s} m_{jk}^{(s_0, s)} - m_{jk}^{(s_0, s)} \right) + \frac{s_0}{1-s} m_n^{(s_0, s)} m_k^{(s_0, s)} m_{\lambda n}^{(s_0, s)} \left( \frac{s_0}{1-s} m_{\lambda n}^{(s_0, s)} - m_{\lambda n}^{(s_0, s)} \right) + \frac{s_0}{1-s} m_n^{(s_0, s)} m_k^{(s_0, s)} m_{\lambda n}^{(s_0, s)} \left( \frac{s_0}{1-s} m_{jk}^{(s_0, s)} - m_{jk}^{(s_0, s)} \right)$$

at  $z = E$ ,

$$\frac{\partial^2 \bar{w}(z; s)}{\partial z^j \partial z^k} = \frac{s_0}{1-s} \left( \frac{s_0}{1-s} m_{\lambda n}^{(s_0, s)} - m_{\lambda n}^{(s_0, s)} \right) \left( \frac{s_0}{1-s} m_{jk}^{(s_0, s)} - m_{jk}^{(s_0, s)} \right) + \frac{s_0}{1-s} m_n^{(s_0, s)} \left( \frac{s_0}{1-s} m_{\lambda n}^{(s_0, s)} - m_{\lambda n}^{(s_0, s)} \right) + \frac{s_0}{1-s} m_n^{(s_0, s)} m_k^{(s_0, s)} m_{\lambda n}^{(s_0, s)} \left( \frac{s_0}{1-s} m_{\lambda n}^{(s_0, s)} - m_{\lambda n}^{(s_0, s)} \right) + \frac{s_0}{1-s} m_n^{(s_0, s)} m_k^{(s_0, s)} m_{\lambda n}^{(s_0, s)} \left( \frac{s_0}{1-s} m_{jk}^{(s_0, s)} - m_{jk}^{(s_0, s)} \right)$$

particular,

The product moments may be obtained from a second differentiation. In

3. Cumulative Processes in  $(t_0, t_0 + t]$

In order to study cumulative processes in an arbitrary interval we define  $W^g(t_0, t)$ , similar to  $W^g(t)$  as given by (2.3.1). Specifically let

$$W^g(t_0, t) = \begin{cases} N(t_0, t) \\ \sum_{n=1}^n g(\bar{J}_{n-1}, J_n, X_n) \\ N(t_0, t) > 0 \end{cases}, \quad N(t_0, t) = 0, \quad (3.3.1)$$

where  $N(t_0, t) = \sum_{j,k} N^{j,k}(t_0, t)$ . Then the stochastic process

$\{W^g(t_0, t); t \geq 0\}$  gives rise to the cumulative process in the arbitrary

interval  $(t_0, t_0 + t]$  for  $t_0$  arbitrary. Define

$$\xi^{j,k}(t_0, p) = \int_0^\infty e^{-pg(j,k,x)} \frac{d}{dx} \phi^{j,k}(t_0, x | \bar{J}(0) = \bar{J}_1) dx \quad (3.3.2)$$

and

$$\xi^{j,k}(s_0, p) = \int_0^\infty e^{-s_0 t_0} \xi^{j,k}(t_0, p) dt_0 \quad (3.3.3)$$

as the moment generating function and L.S.T. of the moment generating

function, respectively, of the first transition after  $t_0$  from state  $j$  to

state  $k$  with sojourn time, measured from  $t_0$ , less than or equal to  $x$ .

Subsequent transitions in  $(t_0, t_0 + t]$  are the same as the transitions

in the interval  $(0, t-x]$  so that the results of Section 2.3 can be em-

ployed. Substituting (3.1.8) into (3.3.2) and applying (3.3.3) gives

$$\xi^{j,k}(s_0, p) = \int_0^\infty e^{-s_0 t_0} \int_0^\infty e^{-pg(j,k,x)} \frac{d}{dx} \phi^{j,k}(t_0, x) dx dt_0 + \int_0^\infty e^{-s_0 t_0} \int_0^\infty e^{-pg(j,k,x)} h_{1j}(t_0 - u) \frac{d}{dx} \phi^{j,k}(u+x) dx dt_0. \quad (3.3.4)$$

The first integral can be written in the form

$$(3.3.5) \quad \int_{-\infty}^0 e^{-s_0 t_0} \theta_{t_0}^{jk}(t_0, p) dt_0,$$

where  $\theta_{t_0}^{jk}(t_0, p)$  is the moment generating function of  $g(j, k, x)$ , when it exists, in the distribution of  $\theta_{t_0}^{jk}(x)$  truncated on the left at  $t_0$ . The L.S.T. with respect to  $t_0$  of (3.3.5) will be denoted by

$$(3.3.6) \quad \theta_{t_0}^{jk}(s_0, p).$$

The second integral in (3.3.4) is

$$\int_{-\infty}^0 e^{-p g(j, k, x)} \frac{d}{dx} \theta_{t_0}^{jk}(u + x) \int_{-\infty}^u e^{-s_0 t_0} \theta_{t_0}^{jk}(t_0 - u) dt_0 du dx.$$

Letting  $v = t_0 - u$  gives

$$\int_{-\infty}^0 e^{-p g(j, k, x)} \frac{d}{dx} \theta_{t_0}^{jk}(u + x) e^{-s_0 v} \int_{-\infty}^0 e^{-s_0 v} \theta_{t_0}^{jk}(v) dv du dx = \int_{-\infty}^0 e^{-p g(j, k, x)} \int_{-\infty}^0 e^{-s_0 u} \frac{d}{dx} \theta_{t_0}^{jk}(u + x) du dx$$

$$(3.3.7) \quad = h_{t_0}^{jk}(s_0, p).$$

So that

$$(3.3.8) \quad \theta_{t_0}^{jk}(s_0, p) = [h_{t_0}^{jk}(s_0, p) + h_{t_0}^{jk}(s_0, p)] \{ [I - p(s_0)]^{-1} \}_{t_0}^{jk}(s_0, p).$$

Since the  $g$ -functions are independent we have

$$(3.3.9) \quad E\{W^g(t_0, t) | N_1^g(t_0, t) = \varepsilon_{jk}(s_0, p)\} = \varepsilon_{jk}(s_0, p) \prod_{m=1}^g \varepsilon_{\delta_{km}}(t_0, t) - \delta_{jk} \delta_{km},$$



$$\begin{aligned}
 & \int_0^{\infty} \sum_{jk} g(j, k, x) \frac{dx}{dt} \bar{J}(0) = e_{-1}^T dx \\
 & = \int_0^{\infty} \sum_{jk} \delta_{jk} g(j, k, x) \frac{dx}{dt} \bar{J}(0) + \int_0^{\infty} \sum_{jk} g(j, k, x) \frac{dx}{dt} \bar{J}(0) + \int_0^{\infty} \sum_{jk} h_{jk}(t_0 - u) \frac{d}{du} A_{jk}(u) \bar{J}(0) du \\
 & \quad + \int_0^{\infty} \sum_{jk} \frac{d}{du} A_{jk}(u) \bar{J}(0) + \int_0^{\infty} \sum_{jk} h_{jk}(t_0 - u) \frac{d}{du} A_{jk}(u) \bar{J}(0) du
 \end{aligned}
 \tag{3.3.12}$$

The expected value of the first term conditional only on  $\bar{J}(0) = e_{-1}^T$ , is

$$\begin{aligned}
 & = g(j, k, x_0) + \sum_{\alpha\beta} N_k^{\alpha\beta}(t - x_0) g(\alpha, \beta, x) \\
 & \quad \{ W^g(t_0, t) | j \rightarrow k \leq x_0, N_k(t - x_0) \}
 \end{aligned}
 \tag{3.3.11}$$

$N_k(t - x_0)$  we write

time, measured from  $t_0$ , less than  $x_0$  and given the transition count matrix

given that the first transition after  $t_0$  is from  $j \rightarrow k$  with sojourn we can still determine its mean and variance.

though we cannot find the L.S.T. of the distribution  $\{ W^g(t_0, t) | \bar{J}(0) = e_{-1}^T \}$

$x_0$ , has a distribution different than subsequent transition times. All difficulty arises because the time to the first transition after  $t_0$ , say  $N_{jk}^g(t_0, t)$  similar to the method used for the interval  $(0, t]$ . The

Unfortunately we are unable to relate this expectation to the p.g.f. of

$$E\{ W^g(t_0, t) | \bar{J}(0) = e_{-1}^T \} = E\{ \sum_{jk} \xi_{jk}(s_0, p) \prod_{m} \xi_{\lambda_m}(p) N_{\lambda_m}^g(t_0, t) - \delta_{jk} \xi_{\lambda_{km}}(p) | \bar{J}(0) = e_{-1}^T \}
 \tag{3.3.10}$$

If the first transition after  $t_0$  is from state  $j$  to state  $k$ . Therefore

Taking expectation with respect to the distribution of  $X$  in the second term

of (3.3.11) gives

$$\int \frac{A^{\alpha\beta}}{P^{\alpha\beta}} M_k^{\alpha\beta}(t - x^0) \cdot$$

Taking expectation with respect to the distribution of the transition count

matrix gives

$$\int \frac{A^{\alpha\beta}}{P^{\alpha\beta}} M_k^{\alpha\beta}(t - x^0)$$

whose expectation with respect to the first transitions  $j \rightarrow k$  is

$$\int_0^t \frac{A^{\alpha\beta}}{P^{\alpha\beta}} \int \delta_{jk} M_k^{\alpha\beta}(t - x) \frac{dx}{d} \bar{J}(0) = \bar{e}_j^T dx$$

$$= \left[ \int_0^t \delta_{jk} M_k^{\alpha\beta}(t - x) \frac{dx}{d} \bar{J}(0) + x | \bar{J}(0) = \bar{e}_j^T dx \right]$$

$$+ \int_0^t \int \delta_{jk} M_k^{\alpha\beta}(t - x) \frac{dx}{d} \bar{J}(0) + x | \bar{J}(0) = \bar{e}_j^T dx$$

$$= \left[ \int_0^t \delta_{jk} M_k^{\alpha\beta}(t - x) \frac{dx}{d} \bar{J}(0) + x | \bar{J}(0) = \bar{e}_j^T dx \right]$$

$$+ \int_0^t \int \delta_{jk} M_k^{\alpha\beta}(t - x) \frac{dx}{d} \bar{J}(0) + x | \bar{J}(0) = \bar{e}_j^T dx$$

(3.3.12)

Combining (3.3.11) and (3.3.12) gives

$$E\{W_g(t_0, \tau) | \bar{J}(0) = \bar{e}_j^T\} = \int_0^t \frac{A^{\alpha\beta}}{P^{\alpha\beta}} M_k^{\alpha\beta}(t_0) \frac{dx}{d} \bar{J}(0) + x | \bar{J}(0) = \bar{e}_j^T dx$$

$$+ \int_0^t \frac{A^{\alpha\beta}}{P^{\alpha\beta}} M_k^{\alpha\beta}(t - x) \frac{dx}{d} \bar{J}(0) + x | \bar{J}(0) = \bar{e}_j^T dx$$

$$E\{W^g(0, t) | \bar{J}(0)\} = \bar{e}_1 \quad (3.3.17)$$

$$E\{W^g(0, t) | \bar{J}(0)\} = \bar{e}_1 = \left[ \sum_k \frac{A_{1k}}{P} + \sum_k \frac{A_{\alpha\beta}}{P} M_{1\alpha\beta}^k(t) - \sum_k \delta_{1\alpha} \frac{A_{\alpha\beta}}{P} \right] + \left[ \sum_k \frac{A_{1k}}{P} + \sum_k \frac{A_{\alpha\beta}}{P} M_{1\alpha\beta}^k(t) - \sum_k \delta_{1\alpha} \frac{A_{\alpha\beta}}{P} \right]$$

Substituting into (3.3.14) gives

$$\int_0^{\infty} \sum_k M_{\alpha\beta}^k(t-x) \frac{d}{dx} \delta_{1k} \alpha(x) dx = M_{1\alpha\beta}^k(t) - \sum_k \delta_{1\alpha} \delta_{k\beta} \quad (3.3.16)$$

that is,

$$M_{1\alpha\beta}^k(t) = \sum_k \delta_{1\alpha} \delta_{k\beta} + \int_0^{\infty} \sum_k M_{\alpha\beta}^k(t-x) \frac{d}{dx} \delta_{1k} \alpha(x) dx$$

in (3.3.15) gives

If the first transition is from state 1 to state k. Taking expectations

$$N_{1\alpha\beta}^k(t) = \delta_{1\alpha} \delta_{k\beta} + N_{\alpha\beta}^k(t-x) \quad (3.3.15)$$

Recall from Section 2.2 that

$$E\{W^g(0, t) | \bar{J}(0)\} = \bar{e}_1 = \left[ \sum_k \frac{A_{1k}}{P} + \sum_k \frac{A_{\alpha\beta}}{P} M_{1\alpha\beta}^k(t) - \sum_k \delta_{1\alpha} \frac{A_{\alpha\beta}}{P} \right] \quad (3.3.14)$$

Note that as  $t_0 \rightarrow 0$

$$\int_0^t \sum_k M_{\alpha\beta}^k(t-x) h_{1j}(t_0-x) \frac{d}{dx} \delta_{jk} \alpha(x) dx \quad (3.3.13)$$

as it should be. Although we have an expression for the mean of  $W^g(t_0, t)$  the reader will note that it involves quantities for which we do not have expressions; namely,  $M_k^{aB}(t-x)$  and  $h_{1j}^{1j}(t_0 - u)$ . However we do have expressions for the L.S.T.'s of these quantities suggesting that we may be able to obtain an expression for the L.S.T. of  $W^g(t_0, t)$  in terms of known L.S.T.'s.

To find the L.S.T. of  $E\{W^g(t_0, t) | \bar{J}(0) = e_{-1}^T\}$ , which we denote by  $w_{1j}^g(s_0, s)$ , first write (3.3.13) in the form

$$\sum_k \frac{A_{1k}(t_0)}{P_{1k}(t_0)} + \int_{t_0}^{\infty} \sum_{j,k} h_{1j}^{1j}(t_0 - u) \frac{A_{jk}(u)}{P_{jk}(u)} du$$

$$+ \sum_{a,B} \frac{A_{aB}^{a,B}}{P_{aB}^{a,B}} \int_{\infty}^0 \sum_{j,k} M_k^{aB}(t-x) \frac{dx}{P_{jk}(t_0)}, \quad x | \bar{J}(0) = e_{-1}^T \quad (3.3.18)$$

Let

$$A_{1k}(s_0) = \int_{-\infty}^0 e^{-s_0 t_0} A_{1k}(t_0) dt_0 \quad (3.3.19)$$

be the L.S.T. of  $A_{1k}(t_0)$ . Then the L.S.T. of (3.3.18) with respect to

$t_0$  is

$$\sum_k \frac{A_{1k}(s_0)}{P_{1k}(s_0)} + \sum_{j,k} \frac{1}{P_{jk}(s_0)} M_{jk}(s_0)$$

$$+ \sum_{a,B} \frac{A_{aB}^{a,B}}{P_{aB}^{a,B}} \int_{-\infty}^0 e^{-s_0 t_0} \sum_{j,k} M_k^{aB}(t-x) \frac{dx}{P_{jk}(t_0)}, \quad x | \bar{J}(0) = e_{-1}^T \quad dx dt_0$$

and taking the L.S.T. with respect to  $t$  gives

$$\frac{1}{s} \sum_{j,k} \frac{A_{jk}(s_0)}{P_{jk}(s_0)} (s_0) + \sum_{a,B} \frac{A_{aB}^{a,B}}{P_{aB}^{a,B}} \int_{-\infty}^0 e^{-s_0 t_0} \sum_{j,k} M_k^{aB}(t-x) \frac{dx}{P_{jk}(t_0)} \quad dx dt_0 \quad (3.3.20)$$

$$\int_{-\infty}^0 e^{-st} M_k^{\alpha\beta}(x-t) \int_{-\infty}^0 e^{-s_0 t_0} \frac{dp}{p} \delta_{jk}^{\alpha\beta}(x+n) \int_{-\infty}^0 e^{-s_0 t_0} h_{\Delta}^{\alpha\beta}(v) dv dx dt =$$

Let  $v = t - n$

$$\int_{-\infty}^0 e^{-st} M_k^{\alpha\beta}(x-t) \int_{-\infty}^0 e^{-s_0 t_0} \frac{dp}{p} \delta_{jk}^{\alpha\beta}(x+n) \int_{-\infty}^0 e^{-s_0 t_0} h_{\Delta}^{\alpha\beta}(t_0 - n) dt_0 dx dt =$$

$$\int_{-\infty}^0 e^{-s_0 t_0} \int_{-\infty}^0 e^{-st} M_k^{\alpha\beta}(x-t) \int_{-\infty}^0 e^{-s_0 t_0} \frac{dp}{p} \delta_{jk}^{\alpha\beta}(x+n) dx dt dt_0$$

$$\frac{s - s_0}{\delta_{jk}^{\alpha\beta}(s)} = \frac{[q_{jk}^{\alpha\beta}(s) - p_{jk}^{\alpha\beta}(s)]}{[r_{jk}^{\alpha\beta}(s)]} \quad (3.3.21)$$

$$\int_{-\infty}^0 \frac{s - s_0}{\delta_{jk}^{\alpha\beta}(s)} e^{-sv} \frac{dv}{p} \delta_{jk}^{\alpha\beta}(v) [1 - e^{-s_0 v}] dv =$$

$$\int_{-\infty}^0 \delta_{jk}^{\alpha\beta}(s) e^{-sv} \frac{dv}{p} \delta_{jk}^{\alpha\beta}(v) \int_{-\infty}^0 e^{-s_0 t_0} dt_0 dv =$$

$$\int_{-\infty}^0 \delta_{jk}^{\alpha\beta}(s) e^{-s_0 t_0} \int_{-\infty}^0 e^{-s(v-t_0)} \frac{dv}{p} \delta_{jk}^{\alpha\beta}(v) dv dt_0$$

Let  $v = t_0 + x$

$$\int_{-\infty}^0 \delta_{jk}^{\alpha\beta}(s) e^{-s_0 t_0} \int_{-\infty}^0 e^{-sx} \frac{dx}{p} \delta_{jk}^{\alpha\beta}(t_0 + x) dx dt_0 =$$

$$\int_{-\infty}^0 \delta_{jk}^{\alpha\beta}(s) e^{-s_0 t_0} \int_{-\infty}^0 \frac{dx}{p} \delta_{jk}^{\alpha\beta}(t_0 + x) \int_{-\infty}^0 e^{-s(v+x)} M_k^{\alpha\beta}(v) dv dx dt_0$$

Let  $v = t - x$

$$\int_{-\infty}^0 \delta_{jk}^{\alpha\beta}(s) e^{-s_0 t_0} \int_{-\infty}^0 \frac{dx}{p} \delta_{jk}^{\alpha\beta}(t_0 + x) \int_{-\infty}^0 e^{-st} M_k^{\alpha\beta}(t - x) dx dt dt_0 =$$

$$\int_{-\infty}^0 \delta_{jk}^{\alpha\beta}(s) e^{-s_0 t_0} \int_{-\infty}^0 \frac{dx}{p} \delta_{jk}^{\alpha\beta}(t_0 + x) \int_{-\infty}^0 e^{-st} M_k^{\alpha\beta}(t - x) dx dt dt_0$$

$$w_{\tau}^{\tau}(s_0, s) = \frac{1}{s} \int \delta_{\tau}^{\tau} \frac{p_{\tau}^{\tau}(s_0)}{m_{\tau}^{\tau}(s_0)} + \left( \int \frac{p_{\tau}^{\tau}(s_0)}{m_{\tau}^{\tau}(s_0)} + \int \frac{p_{\tau}^{\tau}(s_0)}{m_{\tau}^{\tau}(s_0)} \right) + \left\{ \int \frac{p_{\tau}^{\tau}(s_0)}{m_{\tau}^{\tau}(s_0)} - \int \frac{p_{\tau}^{\tau}(s_0)}{m_{\tau}^{\tau}(s_0)} \right\} + \left\{ \int \frac{p_{\tau}^{\tau}(s_0)}{m_{\tau}^{\tau}(s_0)} - \int \frac{p_{\tau}^{\tau}(s_0)}{m_{\tau}^{\tau}(s_0)} \right\}$$

Combining (3.3.20) to (3.3.22) gives

$$\begin{aligned} &= \frac{s - s_0}{m_{\tau}^{\tau}(s_0)} \frac{p_{\tau}^{\tau}(s_0)}{m_{\tau}^{\tau}(s_0)} + \int_0^{\infty} \frac{s - s_0}{m_{\tau}^{\tau}(s_0)} e^{-s v} \left[ -1 \frac{d}{d v} \delta_{\tau}^{\tau}(v) \right] dv \\ &= \int_0^{\infty} \frac{s - s_0}{m_{\tau}^{\tau}(s_0)} e^{-s v} \frac{d}{d v} \delta_{\tau}^{\tau}(v) dv + \int_0^{\infty} e^{-s v} \delta_{\tau}^{\tau}(v) dv \end{aligned}$$

Let  $v = u + x$

$$\begin{aligned} &= \int_0^{\infty} \int_0^{\infty} \frac{s - s_0}{m_{\tau}^{\tau}(s_0)} e^{-s(u+x)} \frac{d}{d x} \delta_{\tau}^{\tau}(u+x) du dx + \int_0^{\infty} \int_0^{\infty} e^{-s(u+x)} \delta_{\tau}^{\tau}(u+x) du dx \\ &= \int_0^{\infty} \int_0^{\infty} \frac{s - s_0}{m_{\tau}^{\tau}(s_0)} e^{-s(u+x)} \frac{d}{d x} \delta_{\tau}^{\tau}(u+x) du dx + \int_0^{\infty} \int_0^{\infty} e^{-s(u+x)} \delta_{\tau}^{\tau}(u+x) du dx \end{aligned}$$

Let  $v = t - x$

$$\begin{aligned} &= \int_0^{\infty} \int_0^{\infty} \frac{s - s_0}{m_{\tau}^{\tau}(s_0)} e^{-s(t-x)} \frac{d}{d x} \delta_{\tau}^{\tau}(t-x) dt dx + \int_0^{\infty} \int_0^{\infty} e^{-s(t-x)} \delta_{\tau}^{\tau}(t-x) dt dx \\ &= \int_0^{\infty} \int_0^{\infty} \frac{s - s_0}{m_{\tau}^{\tau}(s_0)} e^{-s(t-x)} \frac{d}{d x} \delta_{\tau}^{\tau}(t-x) dt dx + \int_0^{\infty} \int_0^{\infty} e^{-s(t-x)} \delta_{\tau}^{\tau}(t-x) dt dx \end{aligned}$$

the g-functions.

which is in terms of the L.S.T.'s of the basic quantities (m, a, g) and

(3.3.24)

$$\begin{aligned} & \left[ \sum_{\alpha} \frac{s_0}{1-s_0} \sum_{\beta} \frac{p_{\alpha\beta}}{A} \left\{ (I - I - p(s)) \right\} \right] \left[ \sum_{\gamma} \frac{s_0}{1-s_0} \sum_{\delta} \frac{q_{\gamma\delta}}{B} \left\{ (I - I - p(s)) \right\} \right] \\ & = \left[ \sum_{\alpha} \frac{s_0}{1-s_0} \sum_{\beta} \frac{p_{\alpha\beta}}{A} \left\{ (I - I - p(s)) \right\} \right] \left[ \sum_{\gamma} \frac{s_0}{1-s_0} \sum_{\delta} \frac{q_{\gamma\delta}}{B} \left\{ (I - I - p(s)) \right\} \right] \\ & \times \left[ \sum_{\alpha} \frac{s_0}{1-s_0} \sum_{\beta} \frac{p_{\alpha\beta}}{A} \left\{ (I - I - p(s)) \right\} \right] \left[ \sum_{\gamma} \frac{s_0}{1-s_0} \sum_{\delta} \frac{q_{\gamma\delta}}{B} \left\{ (I - I - p(s)) \right\} \right] \\ & + \left[ \sum_{\alpha} \frac{s_0}{1-s_0} \sum_{\beta} \frac{p_{\alpha\beta}}{A} \left\{ (I - I - p(s)) \right\} \right] \left[ \sum_{\gamma} \frac{s_0}{1-s_0} \sum_{\delta} \frac{q_{\gamma\delta}}{B} \left\{ (I - I - p(s)) \right\} \right] \\ & = \frac{s_0}{1-s_0} \sum_{\alpha} \frac{p_{\alpha\beta}}{A} \left\{ (I - I - p(s)) \right\} \left[ \sum_{\gamma} \frac{s_0}{1-s_0} \sum_{\delta} \frac{q_{\gamma\delta}}{B} \left\{ (I - I - p(s)) \right\} \right] \end{aligned}$$

So we have

$$m_{\alpha\beta}^k(s) = \sum_{\gamma} \frac{s_0}{1-s_0} \sum_{\delta} \frac{q_{\gamma\delta}}{B} \left\{ (I - I - p(s)) \right\} \left[ \sum_{\epsilon} \frac{s_0}{1-s_0} \sum_{\zeta} \frac{r_{\epsilon\zeta}}{C} \left\{ (I - I - p(s)) \right\} \right]$$

and

$$m_{\alpha\beta}^j(s) = \sum_{\gamma} \frac{s_0}{1-s_0} \sum_{\delta} \frac{q_{\gamma\delta}}{B} \left\{ (I - I - p(s)) \right\} \left[ \sum_{\epsilon} \frac{s_0}{1-s_0} \sum_{\zeta} \frac{r_{\epsilon\zeta}}{C} \left\{ (I - I - p(s)) \right\} \right]$$

used to obtain the alternate expression.

another form. For easier reference we note here those results which are

of known quantities. From previous results we can express (3.3.23) in

So that we have expressed the L.S.T. of  $E\{W^g(t_0, t) | \bar{J}(0) = \bar{e}_1^T\}$  in terms

(3.3.23)

$$\begin{aligned} & \frac{s_0}{1-s_0} \sum_{\alpha} \frac{p_{\alpha\beta}}{A} \left\{ (I - I - p(s)) \right\} \left[ \sum_{\gamma} \frac{s_0}{1-s_0} \sum_{\delta} \frac{q_{\gamma\delta}}{B} \left\{ (I - I - p(s)) \right\} \right] \\ & + \left[ \sum_{\alpha} \frac{s_0}{1-s_0} \sum_{\beta} \frac{p_{\alpha\beta}}{A} \left\{ (I - I - p(s)) \right\} \right] \left[ \sum_{\gamma} \frac{s_0}{1-s_0} \sum_{\delta} \frac{q_{\gamma\delta}}{B} \left\{ (I - I - p(s)) \right\} \right] \\ & = \frac{s_0}{1-s_0} \sum_{\alpha} \frac{p_{\alpha\beta}}{A} \left\{ (I - I - p(s)) \right\} \left[ \sum_{\gamma} \frac{s_0}{1-s_0} \sum_{\delta} \frac{q_{\gamma\delta}}{B} \left\{ (I - I - p(s)) \right\} \right] \end{aligned}$$

To derive the variance of  $W^g(t_0, t | \bar{J}(0) = \bar{e}_1^1)$  we begin with the

well-known identity

$$\text{Var}\{W^g(t_0, t) | \bar{J}(0) = \bar{e}_1^1\} = E \text{Var}\{W^g(t_0, t) | V, \bar{J}(0) = \bar{e}_1^1\} + \text{Var} E\{W^g(t_0, t) | V, \bar{J}(0) = \bar{e}_1^1\} \quad (3.3.25)$$

where  $V$  is the event that the first transition after  $t_0$  is from state  $j$  to state  $k$  with sojourn time, measured from  $t_0$ , less than or equal to  $x$ . Using the same identity for the first term on the right gives

$$\begin{aligned} \text{Var}\{W^g(t_0, t) | \bar{J}(0) = \bar{e}_1^1\} &= E E \text{Var} \left\{ g(j, k, x_0^0) + \int_{t_0}^{t_0+x} N_k^{\alpha\beta} (t - x_0^0) g(\alpha, \beta, x) \right. \\ &+ \left. \text{Var} \left\{ g(j, k, x_0^0) + \int_{t_0}^{t_0+x} M_k^{\alpha\beta} (t - x_0^0) \frac{P}{A} \alpha\beta \right\} \right. \\ &+ \left. E \text{Var} E \left\{ g(j, k, x_0^0) + \int_{t_0}^{t_0+x} N_k^{\alpha\beta} (t - x_0^0) g(\alpha, \beta, x) \right\} \right. \\ &+ \left. \text{Var} \left\{ g(j, k, x_0^0) + \int_{t_0}^{t_0+x} M_k^{\alpha\beta} (t - x_0^0) \frac{P}{A} \alpha\beta \right\} \right. \end{aligned} \quad (3.3.26)$$

For the first term we get

$$E E \int_{t_0}^{t_0+x} N_k^{\alpha\beta} (t - x_0^0) \left[ \frac{B}{A} \alpha\beta - \frac{P}{A} \alpha\beta \right] = E \int_{t_0}^{t_0+x} \left[ \frac{B}{A} \alpha\beta - \frac{P}{A} \alpha\beta \right] M_k^{\alpha\beta} (t - x_0^0)$$

and using (3.3.12) gives for the first term

$$\int_{t_0}^{t_0+x} \left[ \frac{B}{A} \alpha\beta - \frac{P}{A} \alpha\beta \right] M_k^{\alpha\beta} (t - x_0^0) \frac{d}{dx} \int_{t_0}^{t_0+x} N_k^{\alpha\beta} (t - x) \bar{J}(0) = \bar{e}_1^1 dx \quad (3.3.27)$$

The second term is

$$E \text{Var} \left\{ g(j, k, s_0^0) + \int_{t_0}^{t_0+x} \frac{P}{A} \alpha\beta N_k^{\alpha\beta} (t - x_0^0) \right\}$$

$$= \int \int \frac{A}{P} \alpha\beta \gamma \delta \int_{t_0}^{t_0+x} N_k^{\alpha\beta} (t - x) \text{Cov} \left\{ N_k^{\alpha\beta} (t - x), N_k^{\gamma\delta} (t - x) \right\} \frac{d}{dx} \int_{t_0}^{t_0+x} N_k^{\alpha\beta} (t_0, x | \bar{J}(0) = \bar{e}_1^1) dx \quad (3.3.28)$$



The third term is

$$\text{Var} \left\{ g(j, k, x_0) + \sum_{\alpha, \beta} \frac{A}{P} M_{\alpha\beta}^k (t - x_0) \right\}$$

$$= \int_0^{\infty} \sum_{j, k} \left[ g(j, k, x) + \sum_{\alpha, \beta} \frac{A}{P} M_{\alpha\beta}^k (t - x) \right]^2 \frac{dx}{d} \delta_{jk}^k(t_0, x | \bar{J}(0) = \bar{e}_1) dx$$

$$- \left[ \sum_{j, k} \frac{A}{P} M_{jk}^k(t_0) + \sum_{j, k} \int_0^{t_0} h_{1j}^k(t_0 - u) \frac{A}{P} M_{jk}^k(u) du \right]$$

$$+ \sum_{\alpha, \beta} \frac{A}{P} M_{\alpha\beta}^k \int_0^{\infty} \sum_{j, k} M_{\alpha\beta}^k (t - x) \frac{dx}{d} \delta_{jk}^k(t_0, x | \bar{J}(0) = \bar{e}_1) dx \quad (3.3.29)$$

Few simplifications of the last three equations are possible. Since we will ultimately find the L.S.T. of  $\text{Var}\{W^g(t_0, t | \bar{J}(0) = \bar{e}_1)\}$ , we gain little by simplifying the expression for the variance. For easier reference we combine our results giving

$$\text{Var}\{W^g(t_0, t | \bar{J}(0) = \bar{e}_1)\} = \sum_{\alpha, \beta} \left[ \frac{B}{P} M_{\alpha\beta}^k - \frac{A}{P} M_{\alpha\beta}^k \right] \int_0^{\infty} \sum_{j, k} M_{\alpha\beta}^k (t - x) \frac{dx}{d} \delta_{jk}^k(t_0, x | \bar{J}(0) = \bar{e}_1) dx$$

$$+ \sum_{\alpha, \beta} \left[ \sum_{\gamma, \delta} \frac{A}{P} M_{\alpha\beta\gamma\delta}^k + \sum_{\alpha, \beta} \frac{A}{P} M_{\alpha\beta}^k \right] \text{Cov}\{N_k^{\alpha\beta}(t - x), N_k^{\gamma\delta}(t - x)\} \int_0^{\infty} \sum_{j, k} \frac{dx}{d} \delta_{jk}^k(t_0, x | \bar{J}(0) = \bar{e}_1) dx$$

$$+ \int_0^{\infty} \sum_{j, k} \left[ g(j, k, x) + \sum_{\alpha, \beta} \frac{A}{P} M_{\alpha\beta}^k (t - x) \right]^2 \frac{dx}{d} \delta_{jk}^k(t_0, x | \bar{J}(0) = \bar{e}_1) dx$$

$$- \left[ \sum_{j, k} \frac{A}{P} M_{jk}^k(t_0) + \sum_{j, k} \int_0^{t_0} h_{1j}^k(t_0 - u) \frac{A}{P} M_{jk}^k(u) du \right] \int_0^{\infty} \sum_{j, k} \frac{dx}{d} \delta_{jk}^k(t_0, x | \bar{J}(0) = \bar{e}_1) dx \quad (3.3.30)$$

For  $t_0 = 0$ ,

$$\text{Var}\{W^g(0, t) | \bar{J}(0) = e^{-1}\}$$

$$= \int_0^\infty \left[ \frac{B}{P} \frac{\alpha_B}{\alpha_B} - \frac{A}{P} \frac{\alpha_B}{\alpha_B} \right] \left[ \frac{A}{P} \frac{\alpha_B}{\alpha_B} \frac{\gamma_\delta}{\gamma_\delta} \right] \int_0^\infty \text{Cov}\{N_k^{\alpha_B}(t-x), N_k^{\gamma_\delta}(t-x)\} \frac{dx}{dx} \frac{dx}{dx} \\ + \int_0^\infty \left[ g(t, k, x) + \frac{A}{P} \frac{\alpha_B}{\alpha_B} M_k^{\alpha_B}(t-x) \right] \frac{dx}{dx} \frac{dx}{dx}$$

$$- \frac{A}{P} \frac{\alpha_B}{\alpha_B} \frac{\gamma_\delta}{\gamma_\delta} M_k^{\alpha_B}(t) M_k^{\gamma_\delta}(t)$$

(3.3.31)

$$= \left[ \frac{B}{P} \frac{\alpha_B}{\alpha_B} - \frac{A}{P} \frac{\alpha_B}{\alpha_B} \right] \left[ \frac{A}{P} \frac{\alpha_B}{\alpha_B} \frac{\gamma_\delta}{\gamma_\delta} \right] M_k^{\alpha_B}(t) M_k^{\gamma_\delta}(t) - \int_0^\infty \left[ \frac{B}{P} \frac{\alpha_B}{\alpha_B} - \frac{A}{P} \frac{\alpha_B}{\alpha_B} \right] \frac{dx}{dx} \frac{dx}{dx} \\ + \int_0^\infty \left[ \frac{B}{P} \frac{\alpha_B}{\alpha_B} + 2 \frac{A}{P} \frac{\alpha_B}{\alpha_B} \right] \left[ \frac{A}{P} \frac{\alpha_B}{\alpha_B} \frac{\gamma_\delta}{\gamma_\delta} \right] \int_0^\infty \delta_{\alpha_B} \delta_{\gamma_\delta} \frac{dx}{dx} \frac{dx}{dx} \\ - \int_0^\infty \left[ \frac{A}{P} \frac{\alpha_B}{\alpha_B} \frac{\gamma_\delta}{\gamma_\delta} \right] \int_0^\infty \delta_{\alpha_B} \delta_{\gamma_\delta} \frac{dx}{dx} \frac{dx}{dx} \\ + \int_0^\infty \left[ \frac{A}{P} \frac{\alpha_B}{\alpha_B} \frac{\gamma_\delta}{\gamma_\delta} \right] \int_0^\infty \delta_{\alpha_B} \delta_{\gamma_\delta} \frac{dx}{dx} \frac{dx}{dx}$$

$$- \int_{-\infty}^{\infty} \delta_{\alpha\beta}^{\gamma\delta} M_{\beta}^{\gamma\delta}(t-x) \frac{dx}{d} \delta_{\alpha\beta}^{\gamma\delta} \delta_{\alpha\beta}^{\gamma\delta}(x) dx$$

$$- \int_{-\infty}^{\infty} \delta_{\alpha\beta}^{\gamma\delta} M_{\beta}^{\alpha\beta}(t-x) \frac{dx}{d} \delta_{\alpha\beta}^{\gamma\delta} \delta_{\alpha\beta}^{\gamma\delta}(x) dx$$

$$\int_{-\infty}^{\infty} \delta_{\alpha\beta}^{\gamma\delta} M_{\beta}^{\alpha\beta}(t-x) \frac{dx}{d} \delta_{\alpha\beta}^{\gamma\delta} \delta_{\alpha\beta}^{\gamma\delta}(x) dx = \int_{-\infty}^{\infty} \delta_{\alpha\beta}^{\gamma\delta} M_{\beta}^{\alpha\beta}(t-x) \frac{dx}{d} \delta_{\alpha\beta}^{\gamma\delta} \delta_{\alpha\beta}^{\gamma\delta}(x) dx$$

Hence

$$+ \int_{-\infty}^{\infty} \delta_{\alpha\beta}^{\gamma\delta} M_{\beta}^{\alpha\beta}(t-x) \frac{dx}{d} \delta_{\alpha\beta}^{\gamma\delta} \delta_{\alpha\beta}^{\gamma\delta}(x) dx$$

$$+ \int_{-\infty}^{\infty} \delta_{\alpha\beta}^{\gamma\delta} M_{\beta}^{\alpha\beta}(t-x) \frac{dx}{d} \delta_{\alpha\beta}^{\gamma\delta} \delta_{\alpha\beta}^{\gamma\delta}(x) dx$$

$$\int_{-\infty}^{\infty} \delta_{\alpha\beta}^{\gamma\delta} M_{\beta}^{\alpha\beta}(t-x) \frac{dx}{d} \delta_{\alpha\beta}^{\gamma\delta} \delta_{\alpha\beta}^{\gamma\delta}(x) dx = \int_{-\infty}^{\infty} \delta_{\alpha\beta}^{\gamma\delta} M_{\beta}^{\alpha\beta}(t-x) \frac{dx}{d} \delta_{\alpha\beta}^{\gamma\delta} \delta_{\alpha\beta}^{\gamma\delta}(x) dx$$

Taking expectations gives

$$+ \delta_{\alpha\beta}^{\gamma\delta} N_{\beta}^{\alpha\beta}(t-x) + N_{\beta}^{\alpha\beta}(t-x) N_{\beta}^{\alpha\beta}(t-x)$$

$$N_{\beta}^{\alpha\beta}(t-x) = \delta_{\alpha\beta}^{\gamma\delta} N_{\beta}^{\alpha\beta}(t-x) + \delta_{\alpha\beta}^{\gamma\delta} N_{\beta}^{\alpha\beta}(t-x)$$

Now

$$N_{\beta}^{\alpha\beta}(t-x) = E \{ N_{\beta}^{\alpha\beta}(t-x) \}$$

Recall that

$$- \int_{-\infty}^{\infty} \delta_{\alpha\beta}^{\gamma\delta} M_{\beta}^{\alpha\beta}(t-x) \frac{dx}{d} \delta_{\alpha\beta}^{\gamma\delta} \delta_{\alpha\beta}^{\gamma\delta}(x) dx \quad (3.3.32)$$

$$+ \int_{-\infty}^{\infty} \delta_{\alpha\beta}^{\gamma\delta} M_{\beta}^{\alpha\beta}(t-x) \frac{dx}{d} \delta_{\alpha\beta}^{\gamma\delta} \delta_{\alpha\beta}^{\gamma\delta}(x) dx - 2 \int_{-\infty}^{\infty} \delta_{\alpha\beta}^{\gamma\delta} M_{\beta}^{\alpha\beta}(t-x) \frac{dx}{d} \delta_{\alpha\beta}^{\gamma\delta} \delta_{\alpha\beta}^{\gamma\delta}(x) dx$$

but

$$N_1^{\alpha\beta}(t) = \delta_{1\alpha} \delta_{k\beta} + N_k^{\alpha\beta}(t - x)$$

whose expectation, conditional on the first transition, is

$$M_1^{\alpha\beta}(t) = \delta_{1\alpha} \delta_{k\beta} + \int_0^{\infty} \frac{d}{dx} \mathcal{O}_{1k}^{\alpha\beta}(x) dx$$

Using this in the above gives

$$\int_0^{\infty} \mathcal{O}_k^{\alpha\beta, \gamma\delta}(t - x) \frac{d}{dx} \mathcal{O}_{1k}^{\alpha\beta}(x) dx = \mathcal{O}_1^{\alpha\beta, \gamma\delta}(t) - \delta_{1\alpha} \delta_{\beta\delta}$$

$$- \delta_{1\gamma} [M_1^{\alpha\beta}(t) - \delta_{1\alpha} \delta_{\beta\delta}]$$

$$- \delta_{1\gamma} [M_1^{\alpha\beta}(t) - \delta_{1\alpha} \delta_{\beta\delta}]$$

Substituting into (3.3.32) we finally arrive at

$$\text{Var}\{W_g^{\alpha\beta}(0, t) | \bar{J}(0) = \bar{e}_1\} = \left[ \frac{A}{2} \frac{\alpha\beta}{P} - \frac{A}{2} \frac{\alpha\beta}{P} \right] M_1^{\alpha\beta}(t) + \left[ \frac{A}{2} \frac{1k}{P} \right] M_1^{\alpha\beta}(t) + \left[ \frac{A}{2} \frac{1k}{P} \right]$$

$$+ \left[ \frac{A}{2} \frac{\alpha\beta}{P} \right] \left[ \frac{A}{2} \frac{\alpha\beta}{P} \right] \mathcal{O}_1^{\alpha\beta, \gamma\delta}(t) - \delta_{1\alpha} \delta_{\beta\delta} \delta_{1\gamma} \delta_{\delta\delta}$$

$$- \delta_{1\gamma} \delta_{\delta\delta} \delta_{1\alpha} \delta_{\beta\delta} \delta_{1\gamma} + \delta_{1\alpha} \delta_{\beta\delta} \delta_{1\gamma} \delta_{\delta\delta} - \delta_{1\alpha} \delta_{\beta\delta} \delta_{1\gamma} \delta_{\delta\delta} \delta_{1\gamma} \delta_{\delta\delta}$$

$$+ 2 \left[ \frac{A}{2} \frac{\alpha\beta}{P} \right] \left[ \frac{A}{2} \frac{1k}{P} \right] M_1^{\alpha\beta}(t) - 2 \left[ \frac{A}{2} \frac{1k}{P} \right]$$

$$- \left[ \frac{A}{2} \frac{\alpha\beta}{P} \right] \left[ \frac{A}{2} \frac{\alpha\beta}{P} \right] M_1^{\alpha\beta}(t) + \left[ \frac{A}{2} \frac{\alpha\beta}{P} \right] \left[ \frac{A}{2} \frac{\alpha\beta}{P} \right]$$

$$= \left[ \frac{A}{2} \frac{\alpha\beta}{P} - \frac{A}{2} \frac{\alpha\beta}{P} \right] M_1^{\alpha\beta}(t) + \left[ \frac{A}{2} \frac{\alpha\beta}{P} \right] \left[ \frac{A}{2} \frac{\alpha\beta}{P} \right]$$

$$\times \text{COV}\{n_{\alpha\beta}^{\alpha}(t), n_{\gamma\delta}^{\beta}(t)\} + \int \frac{A_2}{P_2} \frac{1}{k} \delta_{\beta\gamma}^{\alpha} \delta_{\delta\alpha}^{\beta}$$

$$- \int \int \frac{A_2}{P_2} \frac{1}{P} \frac{\alpha\beta}{\gamma\delta} M_{\alpha\beta}^{\alpha}(t) + \int \int \frac{A_2}{P_2} \frac{1}{P} \frac{\alpha\beta}{\gamma\delta} \delta_{\beta\gamma}^{\alpha} \delta_{\delta\alpha}^{\beta}$$

$$- \int \int \frac{A_2}{P_2} \frac{1}{P} \frac{\alpha\beta}{\gamma\delta} M_{\alpha\beta}^{\alpha}(t) + 2 \int \int \frac{A_2}{P_2} \frac{1}{P} \frac{\alpha\beta}{\gamma\delta} M_{\alpha\beta}^{\alpha}(t)$$

$$- 2 \int \int \frac{A_2}{P_2} \frac{1}{P} \frac{\alpha\beta}{\gamma\delta}$$

$$= \int \int \frac{A_2}{P_2} \frac{1}{P} \frac{\alpha\beta}{\gamma\delta} - \int \int \frac{A_2}{P_2} \frac{1}{P} \frac{\alpha\beta}{\gamma\delta} M_{\alpha\beta}^{\alpha}(t) + \int \int \frac{A_2}{P_2} \frac{1}{P} \frac{\alpha\beta}{\gamma\delta} \text{COV}\{n_{\alpha\beta}^{\alpha}(t), n_{\gamma\delta}^{\beta}(t)\}$$

$$= \text{Var}\{W^g(t) | \bar{J}(0) = \bar{e}_1\}$$

as required.

We now find the I.S.T. of  $\text{Var}\{W^g(t_0, t) | \bar{J}(0) = \bar{e}_1\}$ . Define

$$\sigma_{\alpha\beta}^{\gamma}(s_0, s) = \int_{-\infty}^0 e^{-s_0 t_0} \int_{-\infty}^0 e^{-st} \text{Var}\{W^g(t_0, t) | \bar{J}(0) = \bar{e}_1\} dt dt_0 \quad (3.3.33)$$

Taking I.S.T.'s term by term in (3.3.30) we have that the first term,

from previous results, is

$$\int_{-\infty}^0 e^{-s_0 t_0} \int_{-\infty}^0 e^{-st} \left[ \int \frac{A_2}{P_2} \frac{\alpha\beta}{\gamma\delta} - \int \frac{A_2}{P_2} \frac{\alpha\beta}{\gamma\delta} M_{\alpha\beta}^{\alpha}(t-x) \frac{dx}{d} \delta_{\beta\gamma}^{\alpha}(t_0, x | \bar{J}(0) = \bar{e}_1) \right] dt dt_0$$

$$= \int \frac{A_2}{P_2} \frac{\alpha\beta}{\gamma\delta} - \int \frac{A_2}{P_2} \frac{\alpha\beta}{\gamma\delta} M_{\alpha\beta}^{\alpha}(s) \delta_{\beta\gamma}^{\alpha}(s_0, s) \quad (3.3.34)$$

For the second term

$$\int \int \frac{A_2}{P_2} \frac{\alpha\beta}{\gamma\delta} \int \int e^{-s_0 t_0} \int_{-\infty}^0 e^{-st} \left[ \int \frac{A_2}{P_2} \frac{\alpha\beta}{\gamma\delta} M_{\alpha\beta}^{\alpha}(t-x) - \int \frac{A_2}{P_2} \frac{\alpha\beta}{\gamma\delta} M_{\alpha\beta}^{\alpha}(t-x) \delta_{\beta\gamma}^{\alpha}(t-x) \right]$$

$$+ \int \int \frac{A_{\alpha\beta} \gamma\delta}{P} \int_0^\infty e^{-st} \int_0^\infty e^{-st} \int_0^\infty e^{-st} M_k^\alpha(t-x) M_k^\gamma(t-x) \frac{d}{dx} Q_{jk}(t_0, x) |j(0) = e_{-1} dx dt dt_0$$

(3.3.36)

Consider the first term only which we expand giving

$$\int_0^\infty \int_0^\infty e^{-st} \int_0^\infty e^{-st} \int_0^\infty e^{-st} \frac{d}{dx} Q_{jk}(t_0 + x) dx dt dt_0$$

$$+ \int_0^\infty \int_0^\infty e^{-st} \int_0^\infty e^{-st} \int_0^\infty e^{-st} \frac{d}{dx} Q_{jk}(t_0 - u) dx dt dt_0$$

$$= \int_0^\infty \int_0^\infty e^{-st} \int_0^\infty e^{-st} \int_0^\infty \frac{d}{dx} Q_{jk}(t_0 + x) dx dt dt_0$$

$$+ \int_0^\infty \int_0^\infty e^{-st} \int_0^\infty \frac{d}{dx} Q_{jk}(u + x) dx dt dt_0 - \int_0^\infty \int_0^\infty e^{-st} \int_0^\infty \frac{d}{dx} Q_{jk}(t_0 - u) dx dt dt_0$$

In the second integral let  $v = t_0 - u$

$$= \int_0^\infty \int_0^\infty e^{-st} \int_0^\infty \frac{d}{dx} Q_{jk}(t_0) dx dt dt_0 + \int_0^\infty \int_0^\infty e^{-st} \int_0^\infty \frac{d}{dx} Q_{jk}(t_0, x) dx dt dt_0$$

$\times Q_{jk}(u+x) dx dt dt_0$

Define

$$\beta_{jk}(s_0) = \int_0^\infty e^{-s_0 t} \int_0^\infty Q_{jk}(t_0) dt dt_0$$

and interchange the order of integration in the second term

$$= \int_0^\infty \int_0^\infty e^{-s_0 u} \int_0^\infty \frac{d}{dx} Q_{jk}(u) dx dt dt_0 + \int_0^\infty \int_0^\infty e^{-s_0 u} \int_0^\infty \frac{d}{dx} Q_{jk}(u+x) dx dt dt_0$$

$$2 \int_0^{\infty} \int_{-\infty}^{\infty} \frac{\alpha_{\beta}}{A} \delta_{\alpha\beta}^{jk} \int_{-\infty}^{\infty} e^{-st} g(j, k, x) M_k^{\alpha\beta}(t-x) \int_{-\infty}^{\infty} \frac{d}{dx} \delta_{\alpha\beta}^{jk}(u+x) e^{-s_0 t} h_{\alpha\beta}^{jk}(t_0-u) dt_0 du dx dt$$

For the second term in write

$$= 2 \int_0^{\infty} \int_{-\infty}^{\infty} \frac{\alpha_{\beta}}{A} \delta_{\alpha\beta}^{jk} \int_{-\infty}^{\infty} \delta_{\alpha\beta}^{jk} (s) \alpha_{jk}(s_0, s) \cdot \quad (3.3.38)$$

Then

$$\alpha_{jk}(s_0, s) = \int_0^{\infty} e^{-s_0 t} \alpha_{jk}(t_0, s) dt_0 \cdot$$

and

$$\alpha_{jk}(t_0, s) = \int_0^{\infty} e^{-sx} g(j, k, x) \frac{d}{dx} \delta_{\alpha\beta}^{jk}(t_0 + x) dx$$

Define

$$= 2 \int_0^{\infty} \int_{-\infty}^{\infty} \frac{\alpha_{\beta}}{A} \delta_{\alpha\beta}^{jk} \int_{-\infty}^{\infty} \delta_{\alpha\beta}^{jk} (s) \int_0^{\infty} e^{-s_0 t} e^{-sx} g(j, k, x) \frac{d}{dx} \delta_{\alpha\beta}^{jk}(t_0+x) dx dt_0 \cdot$$

$$2 \int_0^{\infty} \int_{-\infty}^{\infty} \frac{\alpha_{\beta}}{A} \delta_{\alpha\beta}^{jk} \int_{-\infty}^{\infty} e^{-s_0 t} e^{-sx} g(j, k, x) \frac{d}{dx} \delta_{\alpha\beta}^{jk}(t_0+x) \int_0^{\infty} e^{-sv} M_k^{\alpha\beta}(v) dv dx dt_0$$

Let  $v = t - x$

$$2 \int_0^{\infty} \int_{-\infty}^{\infty} \frac{\alpha_{\beta}}{A} \delta_{\alpha\beta}^{jk} \int_{-\infty}^{\infty} e^{-s_0 t} e^{-sx} g(j, k, x) \frac{d}{dx} \delta_{\alpha\beta}^{jk}(t_0+x) \int_0^x e^{-st} M_k^{\alpha\beta}(t-x) dt dx dt_0$$

For convenience consider the terms individually. The first term becomes

$$+ 2 \int_0^{\infty} \int_{-\infty}^{\infty} \frac{\alpha_{\beta}}{A} \delta_{\alpha\beta}^{jk} \int_{-\infty}^{\infty} e^{-s_0 t} e^{-sx} g(j, k, x) M_k^{\alpha\beta}(t-x) \int_{t_0}^{\infty} h_{\alpha\beta}^{jk}(t_0-u) \frac{d}{dx} \delta_{\alpha\beta}^{jk}(u+x) du dx dt_0 \cdot$$

$$2 \int_0^{\infty} \int_{-\infty}^{\infty} \frac{\alpha_{\beta}}{A} \delta_{\alpha\beta}^{jk} \int_{-\infty}^{\infty} e^{-s_0 t} e^{-sx} g(j, k, x) M_k^{\alpha\beta}(t-x) \frac{d}{dx} \delta_{\alpha\beta}^{jk}(t_0+x) dx dt_0$$

$$\frac{1}{s} \int \frac{\delta_{\alpha\beta}^{\gamma\delta} + m_{\alpha\beta}^{\gamma\delta}(s_0)}{\delta_{\alpha\beta}^{\gamma\delta} + m_{\alpha\beta}^{\gamma\delta}(s_0)} \beta_{jk}^{\alpha\beta}(s_0) + 2 \int \frac{\delta_{\alpha\beta}^{\gamma\delta}}{\delta_{\alpha\beta}^{\gamma\delta} + m_{\alpha\beta}^{\gamma\delta}(s_0)} [\delta_{\alpha\beta}^{\gamma\delta} + m_{\alpha\beta}^{\gamma\delta}(s_0)] m_{\alpha\beta}^{\gamma\delta}(s) \delta_{jk}^{\alpha\beta}(s_0, s)$$

So that the third term in (3.3.30) is

$$(3.3.40) \quad \int \int \frac{\delta_{\alpha\beta}^{\gamma\delta}}{\delta_{\alpha\beta}^{\gamma\delta} + m_{\alpha\beta}^{\gamma\delta}(s_0)} [m_{\alpha\beta}^{\gamma\delta}(s) * m_{\alpha\beta}^{\gamma\delta}(s)] q_{jk}^{\alpha\beta}(s_0, s)$$

which from previous results becomes

$$\int \int \frac{\delta_{\alpha\beta}^{\gamma\delta}}{\delta_{\alpha\beta}^{\gamma\delta} + m_{\alpha\beta}^{\gamma\delta}(s_0)} \int_0^\infty e^{-st} \int_0^\infty e^{-s't} M_{\alpha\beta}^{\gamma\delta}(t-x) M_{\alpha\beta}^{\gamma\delta}(t-x) \frac{dx}{d} \bar{J}(0) = \bar{J}(0) dx dt dt_0$$

The third term in (3.3.36) is

$$(3.3.39) \quad 2 \int \frac{\delta_{\alpha\beta}^{\gamma\delta}}{\delta_{\alpha\beta}^{\gamma\delta} + m_{\alpha\beta}^{\gamma\delta}(s_0)} [\delta_{\alpha\beta}^{\gamma\delta} + m_{\alpha\beta}^{\gamma\delta}(s_0)] m_{\alpha\beta}^{\gamma\delta}(s) \alpha_{jk}^{\alpha\beta}(s_0, s)$$

We have now shown that the second term in (3.3.36) is

$$= 2 \int \frac{\delta_{\alpha\beta}^{\gamma\delta}}{\delta_{\alpha\beta}^{\gamma\delta} + m_{\alpha\beta}^{\gamma\delta}(s_0)} m_{\alpha\beta}^{\gamma\delta}(s_0) \alpha_{jk}^{\alpha\beta}(s_0, s) = 2 \int \frac{\delta_{\alpha\beta}^{\gamma\delta}}{\delta_{\alpha\beta}^{\gamma\delta} + m_{\alpha\beta}^{\gamma\delta}(s_0)} m_{\alpha\beta}^{\gamma\delta}(s_0) \int_0^\infty e^{-s'u} \int_0^\infty e^{-s'x} g(j, k, x) \frac{dx}{d} \bar{J}(0) dx du$$

Let  $v = t - x$

$$= 2 \int \frac{\delta_{\alpha\beta}^{\gamma\delta}}{\delta_{\alpha\beta}^{\gamma\delta} + m_{\alpha\beta}^{\gamma\delta}(s_0)} m_{\alpha\beta}^{\gamma\delta}(s_0) \int_0^\infty g(j, k, x) \int_0^\infty e^{-s_0 u} \frac{dx}{d} \bar{J}(0) dx dt$$

$$2 \int \frac{\delta_{\alpha\beta}^{\gamma\delta}}{\delta_{\alpha\beta}^{\gamma\delta} + m_{\alpha\beta}^{\gamma\delta}(s_0)} m_{\alpha\beta}^{\gamma\delta}(s_0) \int_0^\infty e^{-st} \int_0^\infty g(j, k, x) M_{\alpha\beta}^{\gamma\delta}(t-x) \frac{dx}{d} \bar{J}(0) dx dt$$

Let  $v = t - u$



$$(3.3.41) \quad + \left\{ \frac{A}{\alpha\beta} \frac{\gamma\delta}{P} \right\} [m_k^{\alpha\beta}(s) * m_k^{\gamma\delta}(s)] [p_k^{\alpha\beta}(s_0, s)] \cdot$$

Squaring the last term in (3.3.30) and using (3.3.23) and (3.3.35) we can immediately write that the L.S.T. of the fourth term is

$$\frac{1}{2} \left\{ \frac{h}{P} \right\} \left\{ \frac{h}{P} \right\} \frac{\alpha^{lh}(s_0) * \alpha^{lk}(s_0)}{P^{lh}lk} + \frac{1}{2} \left\{ \frac{h}{P} \right\} \left\{ \frac{h}{P} \right\} \frac{\alpha^{lh}(s_0) * [\alpha^{jk}(s_0) * \alpha^{mj}(s_0)]}{P^{lh}lk}$$

$$+ \frac{1}{2} \left\{ \frac{A}{\alpha\beta} \frac{h}{P} \right\} \left\{ \frac{h}{P} \right\} \frac{\alpha^{lh}(s_0) * [m_k^{\alpha\beta}(s) * p_k^{\alpha\beta}(s_0, s)]}{P^{lh}lk}$$

$$+ \frac{1}{2} \left\{ \frac{A}{\alpha\beta} \frac{h}{P} \right\} \left\{ \frac{h}{P} \right\} \frac{\alpha^{jk}(s_0) * [\alpha^{lm}(s_0) * \alpha^{jm}(s_0)]}{P^{jk}lm}$$

$$+ \frac{1}{2} \left\{ \frac{A}{\alpha\beta} \frac{h}{P} \right\} \left\{ \frac{h}{P} \right\} \frac{\alpha^{jk}(s_0) * [m_k^{\alpha\beta}(s_0) * p_k^{\alpha\beta}(s_0, s)]}{P^{jk}lm}$$

$$+ \left\{ \frac{A}{\alpha\beta} \frac{\gamma\delta}{P} \right\} \left\{ \frac{h}{P} \right\} \left\{ \frac{h}{P} \right\} [m_k^{\alpha\beta}(s) * p_k^{\alpha\beta}(s_0, s)] * [m_k^{\gamma\delta}(s) * p_k^{\gamma\delta}(s_0, s)] \cdot$$

(3.3.42)

Combining the appropriate terms we have shown that the L.S.T. of

$$\text{Var}\{W^g(t_0, t) | \bar{J}(0) = \bar{e}_1\}$$

$$= \left[ \frac{B}{\alpha\beta} - \frac{A}{2} \frac{\alpha\beta}{P} \right] \left\{ \frac{h}{P} \right\} \left\{ \frac{h}{P} \right\} [m_k^{\alpha\beta}(s) * p_k^{\alpha\beta}(s_0, s)] + \left\{ \frac{A}{\alpha\beta} \frac{\gamma\delta}{P} \right\}$$

$$\times \left\{ \frac{C}{\alpha\beta} \frac{\gamma\delta}{P} [m_k^{\alpha\beta}(s) * p_k^{\alpha\beta}(s_0, s)] - [m_k^{\alpha\beta}(s) * m_k^{\gamma\delta}(s)] [p_k^{\alpha\beta}(s_0, s)] \right\}$$

$$+ \frac{1}{2} \left\{ \frac{h}{P} \right\} \left\{ \frac{h}{P} \right\} \frac{\alpha^{jk}(s_0) * [m_k^{\alpha\beta}(s_0) * p_k^{\alpha\beta}(s_0, s)]}{P^{jk}lm} + 2 \left\{ \frac{A}{\alpha\beta} \frac{h}{P} \right\} \left\{ \frac{h}{P} \right\} \frac{\alpha^{jk}(s_0) * [m_k^{\alpha\beta}(s_0) * p_k^{\alpha\beta}(s_0, s)]}{P^{jk}lm}$$

$$+ \left\{ \frac{A}{\alpha\beta} \frac{\gamma\delta}{P} \right\} \left\{ \frac{h}{P} \right\} \left\{ \frac{h}{P} \right\} [m_k^{\alpha\beta}(s) * m_k^{\gamma\delta}(s)] [p_k^{\alpha\beta}(s_0, s)] + \left\{ \frac{A}{\alpha\beta} \frac{\gamma\delta}{P} \right\}$$

$$- \frac{1}{2} \left\{ \frac{h}{P} \right\} \left\{ \frac{h}{P} \right\} \frac{\alpha^{lh}(s_0) * \alpha^{lk}(s_0)}{P^{lh}lk}$$

$$-\frac{2}{s} \sum_h \sum_k \frac{1}{P} \alpha^{jh} (s_0) * [\alpha^{jk} (s_0) m_{1j}^k (s_0)]$$

$$-\frac{2}{s} \sum_h \sum_k \frac{1}{P} \alpha^{jh} (s_0) * [m_k^{\alpha\beta} (s_0) q_k^{*j} (s_0), s]$$

$$-\frac{1}{s} \sum_k \sum_m \frac{1}{P} \alpha^{jk} (s_0) m_{1j}^k (s_0) * [\alpha^{\lambda m} (s_0) m_{1j}^{\lambda} (s_0)]$$

$$-\frac{2}{s} \sum_k \sum_m \frac{1}{P} \alpha^{jk} (s_0) m_{1j}^k (s_0) * [m_m^{\alpha\beta} (s_0) q_m^{*j} (s_0), s]$$

$$-\sum_k \sum_m \frac{1}{P} \alpha^{jk} (s_0) m_{1j}^k (s_0) * [m_m^{\lambda\delta} (s_0) q_m^{*j} (s_0), s]$$

(3.3.43)

4. Asymptotic Results

Recall that the L.S.T. of  $E\{N_{1j}^k(t_0, t)\}$  is the  $i$ <sup>th</sup> element of the

vector  $m_{1j}^k(s_0, s)$ ; namely

$$m_{1j}^k(s_0, s) = \frac{s_0}{1-s} (m_{1j}^k(s_0) - m_{1j}^k(s_0)) \quad (3.4.1)$$

For small  $s$ ; i.e., large  $t$ , replace  $m_{1j}^k(s)$  with (2.4.3). Then

$$m_{1j}^k(s_0, s) = \frac{s_0}{1-s} [P^{jk} \left\{ \frac{\alpha}{s} v_j + av_j + \alpha(h_{1j} - v_j) \right\} + s \left( \delta_{1j}^{jk} - \mu^{jk} \{av_j + \alpha h_{1j}\} + \frac{2}{1} \alpha v_j \mu^{jk} \right) + o(s)]$$

$$= \frac{s_0}{1} \left[ \frac{s}{\alpha v_j P^{jk}} + av_j P^{jk} + \alpha P^{jk} (h_{1j} - v_j) - m_{1j}^k(s_0) \right] + \frac{s_0}{\alpha v_j P^{jk}} + \frac{s_0}{av_j P^{jk}} + \frac{s_0^2}{2 P^{jk}} \left\{ \frac{s_0}{\alpha v_j P^{jk} (h_{1j} - v_j)} + \frac{s_0}{av_j P^{jk}} + \frac{s_0}{2 P^{jk}} \right\}$$

$$- \frac{s_0}{m_{1j}^k(s_0)} + P^{jk} \left( \delta_{1j}^{jk} - \mu^{jk} \{av_j + \alpha h_{1j}\} \right)$$

$$+ \frac{2}{1} \alpha v_j \mu^{jk} \left( \right) + o(s)$$

(3.4.2)

whose inverse with respect to  $s$  is, for large  $t$ ,

$$m_{jk}^{\frac{1}{t}}(s_0, t) = \frac{1}{s_0} \left[ \alpha v_{j,p}^j t + \alpha v_{j,p}^j k + \alpha p_{jk}^j (h_{1j} - v_{j,\mu}^j k) \right]$$

$$- m_{jk}^{\frac{1}{t}}(s_0) + \frac{s_0}{\alpha v_{j,p}^j k} + o(1) \left[ \right]$$

and inverting with respect to  $s_0$  gives

$$M_{jk}^{\frac{1}{t}}(t_0, t) = \alpha v_{j,p}^j t_0 + \alpha v_{j,p}^j k t + \alpha v_{j,p}^j k + \alpha p_{jk}^j (h_{1j} - v_{j,\mu}^j k) - M_{jk}^{\frac{1}{t}}(t_0) + o(1) \quad (3.4.3)$$

Note that for  $t_0 = 0$ ,  $M_{jk}^{\frac{1}{t}}(0, t) = M_{jk}^{\frac{1}{t}}(t)$  for large  $t$ . Also

$$\lim_{t \rightarrow \infty} \frac{t}{M_{jk}^{\frac{1}{t}}(t_0, t)} = \alpha v_{j,p}^j k \quad (3.4.4)$$

which is independent of both  $t_0$  and the initial state. It can be further

shown, for  $j \neq \lambda$  or  $k \neq n$ , that

$$c_{jk,\lambda n}^{\frac{1}{t}}(t_0, t) = \alpha^2 p_{j,p}^j k v_{j,\mu}^j t^2 + \left[ \alpha v_{j,p}^j k \{ \alpha v_{j,p}^j k + \alpha p_{\lambda n}^j (h_{1\lambda} - v_{j,\mu}^j k) \} + \alpha p_{\lambda n}^j (h_{1\lambda} - v_{j,\mu}^j k) \right]$$

$$- M_{\lambda n}^{\frac{1}{t}}(t_0) + \alpha v_{j,p}^j k t_0 + p_{jk}^j \{ \alpha v_{j,p}^j k + \alpha (h_{nj} - v_{j,\mu}^j k) \} \{ \alpha v_{j,p}^j k \}$$

$$+ \alpha v_{j,p}^j k \{ \alpha v_{j,p}^j k + \alpha p_{jk}^j (h_{1j} - v_{j,\mu}^j k) - M_{jk}^{\frac{1}{t}}(t_0) \} + \alpha v_{j,p}^j k t_0$$

$$+ p_{\lambda n}^j \{ \alpha v_{j,p}^j k + \alpha (h_{k\lambda} - v_{j,\mu}^j k) \} \{ \alpha v_{j,p}^j k \} t + o(t) \quad (3.4.5)$$

For  $t_0 = 0$  this reduces to  $c_{jk,\lambda n}^{\frac{1}{t}}(t)$  as required. We also can show that

$$\text{Cov} \{ N_{jk}^{\frac{1}{t}}(t_0, t), N_{\lambda n}^{\frac{1}{t}}(t_0, t) \} = 2 \alpha v_{j,p}^j k \{ \alpha v_{j,p}^j k + \alpha (h_{nj} - v_{j,\mu}^j k) \}$$

$$+ \alpha^2 v_{j,p}^j k p_{jk}^j k (h_{k\lambda} - v_{j,\mu}^j k) - \alpha^2 v_{j,p}^j k p_{jk}^j k (h_{nj} - v_{j,\mu}^j k)$$

$$- 2 \alpha v_{j,p}^j k \{ \alpha v_{j,p}^j k + \alpha (h_{1\lambda} - v_{j,\mu}^j k) \} - \alpha^2 v_{j,p}^j k p_{jk}^j k (h_{1\lambda} - v_{j,\mu}^j k)$$

$$- \alpha^2 v_{j,p}^j k p_{jk}^j k (h_{1j} - v_{j,\mu}^j k) + v_{j,p}^j k M_{\lambda n}^{\frac{1}{t}}(t_0) (\alpha t_0 + a)$$

$$\begin{aligned}
 & + v_{j,p}^j \lambda_{n,jk}^j M_{jk}^j(t_0) (\alpha t_0 + a) + \alpha p_{jk}^j (h_{1j}^j - v_{j,\mu}^j) M_{jk}^j(t_0) \\
 & + v_{j,p}^j \lambda_{n,jk}^j M_{jk}^j(t_0) (\alpha t_0 + a) + \alpha p_{jk}^j (h_{1j}^j - v_{j,\mu}^j) M_{jk}^j(t_0) + o(t)
 \end{aligned}
 \tag{3.4.6}$$

and

$$\text{Var}\{N_{jk}^j(t_0, t)\} = \alpha v_{j,p}^j t + 2\alpha v_{j,p}^j t^2 + \alpha^2 v_{j,p}^j t^3 - \alpha^2 v_{j,p}^j t^2 (t - t_0) - \alpha^2 v_{j,p}^j t^2 t_0$$

$$\begin{aligned}
 & + 2\alpha^2 v_{j,p}^j t (h_{kj}^j - v_{j,\mu}^j) - 2\alpha^2 v_{j,p}^j t_0 (h_{1j}^j - v_{j,\mu}^j) \\
 & + 2v_{j,p}^j M_{jk}^j(t_0) (\alpha t_0 + a) + 2\alpha p_{jk}^j (h_{1j}^j - v_{j,\mu}^j) M_{jk}^j(t_0)
 \end{aligned}$$

$$- \left( M_{jk}^j(t_0) \right)^2 + o(t) .
 \tag{3.4.7}$$

Furthermore

$$\lim_{t \rightarrow \infty} \frac{1}{t} \text{Cov}\{N_{jk}^j(t_0, t), N_{jn}^j(t_0, t)\} =$$

$$= 2\alpha v_{j,p}^j v_{j,p}^j \lambda_{n,jk}^j + \alpha^2 v_{j,p}^j \lambda_{n,jk}^j (h_{nj}^j - v_{j,\mu}^j) + \alpha^2 v_{j,p}^j \lambda_{n,jk}^j (h_{kj}^j - v_{j,\mu}^j) - v_{j,\mu}^j \lambda_{n,jk}^j$$

and

$$\lim_{t \rightarrow \infty} \frac{1}{t} \text{Var}\{N_{jk}^j(t_0, t)\} = \alpha v_{j,p}^j + 2\alpha v_{j,p}^j t + 2\alpha^2 v_{j,p}^j t^2 + 2\alpha^2 v_{j,p}^j t_0 (h_{kj}^j - v_{j,\mu}^j)
 \tag{3.4.9}$$

Lemma:  $(I + pQ)^{-1} = I - p(I + pQ)^{-1}Q$ .

(4.1.2)  $(I - q(s))^{-1} = (I - p \bar{F}'(s))^{-1}$ .

of L.S.P.'s of the  $F_i(t)$ . Then

$Q_{ij}^{(1)}(t) = p_{ij}^{(1)}(t)$ , hence  $q(s) = p \bar{F}'(s)$ , where  $\bar{F}'(s)$  is the column vector

where  $\bar{p}'$  is the row vector of the  $P_i$ . Also the matrix  $Q(t)$  has  $(i, j)$ th element

(4.1.1)  $P_0 = \begin{bmatrix} P_1 & P_2 & \dots & P_m \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ P_1' & P_2' & \dots & P_m' \\ P_1 & P_2 & \dots & P_m \end{bmatrix} = \bar{e} \bar{p}'$ ,

of the form

For this special case the matrix  $P_0$  of transition probabilities is

1. Zero-Order M.R.P.'s

to the case of  $m$  states.

special case we treat the two-state M.R.P. and derive results analogous

$F_{ij}^{(1)}(x)$  is a function of either  $i$  or  $j$  but not both. For the second

Order M.R.P.'s. Recall from (1.1) that a Zero-Order M.R.P. arises when

and derive the results for two special cases. First we consider zero-

In this chapter we specialize the results of the preceding chapters

SPECIAL CASES

CHAPTER IV

$$(4.1.6) \quad \left\{ (I - \bar{p}\bar{f}'(s))^{-1} \right\}_{1j} = \delta_{1j} + p_j \left[ \frac{1}{s} + \frac{1}{2} \gamma_2 - \mu_1 \right] + \frac{1}{2} \gamma_3 + \mu_1 \left[ \frac{1}{s} + \frac{1}{2} \gamma_2 - \mu_1 \right] + o(s) \quad ,$$

In powers of  $s$  then (4.1.4) is and the  $\mu_k^{(1)}$  are the  $1^{\text{th}}$  order raw moments of the  $F_k(x)$ ,  $k = 1, 2, \dots, m$ .

$$\mu^{(1)} = \sum_{k=1}^m p_k \mu_k^{(1)}$$

where

$$(4.1.5) \quad \left\{ p_k F_k(s) \right\} = 1 - s \mu_1 + \frac{1}{2} s^2 \mu_2 - \frac{1}{6} s^3 \mu_3 + o(s^3) \quad ,$$

Let

$$(4.1.4) \quad \left\{ (I - \bar{p}\bar{f}'(s))^{-1} \right\}_{1j} = \delta_{1j} + p_j q_1(s) \left[ 1 - \sum_{k=1}^k p_k F_k(s) \right]^{-1} \quad .$$

Hence for the zero-order M.R.P. the basic quantity  $(I - q(s))^{-1}$  can be exactly determined. From (4.1.3) we immediately have

$$(4.1.3) \quad \bar{p}\bar{f}'(s) = I + \frac{1 - \bar{f}'(s)\bar{p}}{\bar{p}\bar{f}'(s)} \quad .$$

$$(I - \bar{p}\bar{f}'(s))^{-1} = I + \bar{p} \left( I - \bar{f}'(s)\bar{p} \right)^{-1} \bar{p}\bar{f}'(s)$$

Using the Lemma (4.1.2) becomes

$$\begin{aligned} &= I - p(I + \bar{p}\bar{f}'(s))^{-1} \bar{p} \quad . \\ &= I - p[I - \bar{p}\bar{f}'(s)]^{-1} \bar{p} + (\bar{p}\bar{f}'(s))^2 - \dots \\ \text{Proof: } & (I + p\bar{q})^{-1} = I - p\bar{q} + (p\bar{q})^2 - (p\bar{q})^3 + \dots \end{aligned}$$

where

$$(4.1.7) \quad \gamma_2 = \frac{n \cdot}{n} \quad , \quad \gamma_3 = \frac{n \cdot}{n} \quad (3)$$

Using (4.1.6) in (2.2.5) we get

$$m_{jk}^i(s) = p_{F_j}^k(s) \{ (I - \bar{p}F^i(s))^{-1} \}_{i,j}$$

$$= \frac{1}{s} p_{F_j}^k + p_{F_j}^k \delta_{ij} + p_{F_j}^k \left( \frac{1}{s} \gamma_2 - \gamma_1 - n_j \right)$$

$$+ s \left\{ p_{F_j}^k \left[ \frac{1}{s} \gamma_2 + n_j \right] + \left( \frac{1}{s} \gamma_2 - n_j \right) \left( \frac{1}{s} \gamma_2 - n_j \right) - \frac{1}{s} \gamma_3 \right\}$$

$$(4.1.8) \quad - \delta_{ij} p_{F_j}^k + o(s) \cdot$$

Formal inversion gives, for large  $t$ ,

$$(4.1.9) \quad m_{jk}^i(t) = p_{F_j}^k t + p_{F_j}^k \delta_{ij} + p_{F_j}^k \left( \frac{1}{t} \gamma_2 - n_j \right) + o(1) \cdot$$

Using (4.1.8) and the results of Section 2.2 we obtain, after a little

algebra,

$$\text{Cov} \{ N_{jk}^i(t), N_{kn}^j(t) \} = \delta_{jk} \delta_{kn} p_{F_j}^k t + p_{F_j}^k \left( \delta_{kn} \delta_{ij} + \delta_{nj} p_{F_j}^k \right)$$

$$+ p_{F_j}^k p_{F_j}^k \delta_{kn} \left( \frac{1}{t} \gamma_2 - n_j - n_k - n_j - n_k - n_j \right)$$

$$(4.1.10) \quad + o(t)$$

and in particular

$$(4.1.11) \quad \text{Var} \{ N_{jk}^i(t) \} = p_{F_j}^k t + \frac{n \cdot}{2 p_{F_j}^k} \delta_{jk} t + \left( p_{F_j}^k \right)^2 \left( \frac{1}{t} \gamma_2 - 2 n_j - 2 n_k - 2 n_j \right) + o(t)$$

For large  $t$  we have shown for zero-order M.R.P.'s that

$$(4.1.12) \quad \lim_{t \rightarrow \infty} \frac{1}{t} M_{jk}^I(t) = \frac{P_{jk}^I}{n},$$

$$\lim_{t \rightarrow \infty} \frac{1}{t} \text{COV}\{N_{jk}^I(t), N_{kn}^I(t)\} = \delta_{jk} \delta_{kn} \frac{P_{jk}^I}{n} + \frac{P_{kn}^I}{n} (\delta_{kj} \delta_{kn} + \delta_{nj} \delta_{k})$$

$$(4.1.13) \quad + \frac{P_{jk}^I P_{kn}^I}{n} (\gamma_2 - \gamma_j - \gamma_k - \gamma_n),$$

and

$$(4.1.14) \quad \lim_{t \rightarrow \infty} \frac{1}{t} \text{Var}\{N_{jk}^I(t)\} = \frac{P_{jk}^I}{n} + \frac{2P_{jk}^I}{n} \delta_{jk} + \left(\frac{P_{jk}^I}{n}\right)^2 (\gamma_2 - 2\gamma_j - 2\gamma_k).$$

Note that these asymptotic moments are independent of the initial state

of the process, as expected.

### 2. Two-State M.R.P.'s

For the special case  $m = 2$  we write

$$(I-q(s)) = \begin{bmatrix} 1 - a + sc_{11} - \frac{1}{2} s^2 d_{11} + \frac{1}{6} s^3 e_{11} & a - 1 + sc_{12} - \frac{1}{2} s^2 d_{12} + \frac{1}{6} s^3 e_{12} \\ -b + sc_{21} - \frac{1}{2} s^2 d_{21} + \frac{1}{6} s^3 e_{21} & b + sc_{22} - \frac{1}{2} s^2 d_{22} + \frac{1}{6} s^3 e_{22} \end{bmatrix}$$

where

$$(4.2.1) \quad P_{11} = 1 - P_{12} = a, \quad P_{21} = 1 - P_{22} = b.$$

Then

$$(4.2.2) \quad \det(I - q(s)) = as + bs^2 + \gamma s^3 + o(s^3),$$

where

$$\alpha = (1 - a)(c_{21} + c_{22}) + b(c_{11} + c_{12}),$$



we may obtain, by direct evaluation, the means for the two-state process.

$$(4.2.4) \quad = q_{jk}^{-1}(s) a_{1j}(s)$$

$$m_{jk}^{-1}(s) = q_{jk}^{-1}(s) \left( I - q(s) \right)^{-1}$$

From the previous result for m-state processes

$$\theta_1 = \beta_2 - \alpha\gamma$$

where

$$a_{22}(s) = \frac{1}{s} \left( \frac{\alpha}{b} \right) + \frac{1}{2} [\alpha c_{22} - b\beta] + \frac{\alpha}{s} [\theta_1 b - \alpha\beta c_{22} - \frac{1}{2} \alpha^2 d_{22}] + o(s)$$

and

$$a_{21}(s) = \frac{1}{s} \left( \frac{\alpha}{1-a} \right) - \frac{1}{2} [\alpha c_{12} + \beta(1-a)] + \frac{\alpha}{s} [\theta_1(1-a) + \alpha\beta c_{12} + \frac{1}{2} \alpha^2 d_{12}] + o(s)$$

$$a_{12}(s) = \frac{1}{s} \left( \frac{\alpha}{b} \right) - \frac{1}{2} [\alpha c_{21} + b\beta] + \frac{\alpha}{s} [\theta_1 b + \alpha\beta c_{21} + \frac{1}{2} \alpha^2 d_{21}] + o(s)$$

$$a_{11}(s) = \frac{1}{s} \left( \frac{\alpha}{1-a} \right) + \frac{1}{2} [\alpha c_{11} - \beta(1-a)] + \frac{\alpha}{s} [\theta_1(1-a) - \alpha\beta c_{11} - \frac{1}{2} \alpha^2 d_{11}] + o(s)$$

Letting  $A(s) = (I - q(s))^{-1}$  we get, omitting the algebra,

$$(4.2.3) \quad \det(I - q(s)) = \frac{1}{1 - \frac{\alpha s}{\beta} - \frac{\alpha^2}{\beta} + \frac{\alpha^3}{s} (\beta_2 - \alpha\gamma) + o(s)}$$

so that

$$P_1 = \begin{bmatrix} p_{11} & p_{12} \\ p_{21} & p_{22} \end{bmatrix}$$

$$\gamma = \frac{1}{6} \left\{ (1-a)(e_{21} + e_{22}) + b(e_{11} + e_{12}) - \frac{1}{2} [c_{11}d_{22} + c_{22}d_{11} - c_{12}d_{21} - c_{21}d_{12}] \right\}$$

$$\beta = \det P_1 - \frac{1}{2} \left\{ (1-a)(d_{21} + d_{22}) + b(d_{11} + d_{12}) \right\}$$

We omit the algebraic details, but one can easily show that

$$P_{11}^{m_{11}}(s) = \frac{P_{11}^{\alpha_3}}{1} \left\{ \frac{1}{2} \alpha^2 (1 - a) + \alpha^2 c_{11} - \alpha (1 - a) (\beta + \alpha \mu_{11}) \right. \\ \left. + s[\theta_1 (1 - a) - \alpha \beta c_{11} - \frac{1}{2} \alpha^2 d_{11} - \mu_{11} (\alpha^2 c_{11} - \alpha \beta (1 - a))] \right. \\ \left. + \frac{1}{2} \mu_{11} \alpha^2 (1 - a) \right] + o(s) \} \quad (4.2.5)$$

To compute  $m_{12}^{(k)}$  replace  $\mu_{11}^{(k)}$  by  $\mu_{12}$  and  $P_{11}$  by  $P_{12}$  in (4.2.5).

$$P_{22}^{m_{22}}(s) = \frac{P_{22}^{\alpha_3}}{1} \left\{ \frac{1}{2} \alpha^2 b - \alpha^2 c_{21} - \alpha b (\beta + \alpha \mu_{22}) \right. \\ \left. + s[\theta_1 b + \alpha \beta c_{21} + \frac{1}{2} \alpha^2 d_{21} + \mu_{22} (\alpha^2 c_{21} + \alpha b (\beta + \alpha \mu_{22})) \right. \\ \left. + \frac{1}{2} \mu_{22} \alpha^2 b \right] + o(s) \} \quad (4.2.6)$$

To compute  $m_{21}^{(k)}$  replace  $\mu_{22}^{(k)}$  by  $\mu_{21}$  and  $P_{22}$  by  $P_{21}$  in (4.2.6).

$$P_{11}^{m_{11}}(s) = \frac{P_{11}^{\alpha_3}}{1} \left\{ \frac{1}{2} \alpha^2 (1 - a) - \alpha^2 c_{12} - \alpha (1 - a) (\beta + \alpha \mu_{11}) \right. \\ \left. + s[\theta_1 (1 - a) + \alpha \beta c_{12} + \frac{1}{2} \alpha^2 d_{12} + \mu_{11} (\alpha^2 c_{12} + \alpha \beta (1 - a))] \right. \\ \left. + \frac{1}{2} \mu_{11} \alpha^2 (1 - a) \right] + o(s) \} \quad (4.2.7)$$

To compute  $m_{12}^{(k)}$  replace  $\mu_{11}^{(k)}$  by  $\mu_{12}$  and  $P_{11}$  by  $P_{12}$  in (4.2.7).

$$P_{22}^{m_{22}}(s) = \frac{P_{22}^{\alpha_3}}{1} \left\{ \frac{1}{2} \alpha^2 b + \alpha^2 c_{22} - \alpha b (\beta + \alpha \mu_{22}) \right. \\ \left. + s[\theta_1 - \alpha \beta c_{22} - \frac{1}{2} \alpha^2 d_{22} - \mu_{22} (\alpha^2 c_{22} - \alpha \beta b) \right. \\ \left. + \frac{1}{2} \mu_{22} \alpha^2 b \right] + o(s) \} \quad (4.2.8)$$

To compute  $m_{21}^{(k)}$  replace  $\mu_{22}^{(k)}$  by  $\mu_{21}$  and  $P_{22}$  by  $P_{21}$  in (4.2.8). From

the above expressions we obtain

$$\lim_{t \rightarrow \infty} \frac{1}{t} M_{11}^t(t) = \frac{P_{11}^{\alpha_3}}{1} (1 - a) = \frac{\alpha}{1} P_{11} (1 - a) \quad (4.2.9)$$

$$\lim_{t \rightarrow \infty} \frac{1}{t} M_1(t) = \frac{p_{22}}{\alpha} b = \frac{\alpha}{1} p_{22} (1 - p_{22}), \quad (4.2.10)$$

$$\lim_{t \rightarrow \infty} \frac{1}{t} M_2(t) = \frac{p_{11}}{\alpha} (1 - a) = \frac{\alpha}{1} p_{11} (1 - p_{11}), \quad (4.2.11)$$

$$\lim_{t \rightarrow \infty} \frac{1}{t} M_2(t) = \frac{\alpha}{p_{22}} b = \frac{\alpha}{1} p_{22} (1 - p_{22}), \quad (4.2.12)$$

recall that

$$\alpha = (1 - a)(c_{21} + c_{22}) + b(c_{11} + c_{12})$$

$$= (1 - p_{11})(p_{21}^{11} + p_{22}^{11}) + (1 - p_{22})(p_{11}^{11} + p_{12}^{11}), \quad (4.2.13)$$

We may summarize and write in matrix notation

$$\lim_{t \rightarrow \infty} \frac{1}{t} M_1(t) = \lim_{t \rightarrow \infty} \frac{1}{t} M_2(t) = \frac{\alpha}{1} \begin{bmatrix} p_{11} p_{12} & p_{22} p_{21} \\ p_{11} p_{12} & p_{22} p_{21} \end{bmatrix}. \quad (4.2.14)$$

$$\lim_{s \rightarrow 0} \frac{1}{s} |I - P_0 + sP_1| = \lim_{s \rightarrow 0} \frac{d}{ds} |I - P_0 + sP_1| \cdot$$

But  $\lim_{s \rightarrow 0} |I - P_0 + sP_1| = |I - P_0| = 0$ , hence using L'Hopital's Rule

$$\frac{1}{\alpha} = \det(P_1) \prod_{i=1}^{m-1} \lambda_i = \lim_{s \rightarrow 0} \frac{1}{s} |I - P_0 + sP_1| \cdot$$

Now, from (1.7.3)

$$\eta_i = \prod_{j=1}^j P_{ij} \eta_j \cdot \quad (5.3)$$

where

$$\frac{1}{\alpha} = \bar{v}' \bar{\eta} \cdot \quad (5.2)$$

that

where the  $\lambda_i$  are the nonzero latent roots of  $P_1^{-1}(I - P_0)$ . We shall show

$$\alpha = \det(P_1) \prod_{i=1}^{m-1} \lambda_i \cdot \quad (5.1)$$

First recall from Section 1.7 that

rather than in terms of the mean recurrence times as in Pyke (1961b).

are derived. Results are given in terms of the basic quantities  $\tilde{Q}(t)$

In this chapter certain limiting distributions for previous results

## LIMITING DISTRIBUTIONS

### CHAPTER V

$$P_{ij}(t) = \Pr\{J(t) = \bar{j} | J(0) = \bar{i}\} \quad (5.5)$$

Before discussing the limiting results for the distributions derived in this paper we will use (5.4) to give alternate limiting expressions for results due to Pyke (1961b). First, define

$$\frac{1}{s} \frac{d}{ds} (I - P(s))^{-1} = \frac{1}{s} \frac{dV}{dV} + \text{constant} + o(1) \quad (5.4)$$

results in hand we can write

abilities for the transition probability matrix  $P_0$ . With the above vector for  $\lambda = 1$  in  $|P_0 - \lambda I| = 0$  and therefore are the steady-state probabilities we have shown that  $\frac{\alpha}{1} = \bar{v}' \bar{\eta}$ . Note also that  $\bar{v}'$  is the left eigen-

$$\lim_{s \rightarrow 0} \frac{d}{ds} |I - P_0 + sP_1| = \sum_{i,j} \bar{v}'_i P_{ij} \bar{\eta}_j = \bar{v}' P_1 \bar{\eta} = \bar{v}' \bar{\eta} \quad .$$

any row of  $(I - P_0)$  are the same. Let  $v_i$  be these cofactors. Then but  $(I - P_0)$  is singular and therefore the cofactors of the elements in

$$\lim_{s \rightarrow 0} \frac{d}{ds} |I - P_0 + sP_1| = \sum_{i,j} P_{ij} v_i \bar{\eta}_j = \sum_{i,j} P_{ij} v_i \bar{\eta}_j |I - P_0|_{ij} \quad ,$$

Hence  $|A + sB|$

where  $B = (b_{ij})$  and  $|A + sB|_{ij}$  is the cofactor of the  $(i,j)$ th element of

$$\frac{d}{ds} |A + sB| = \sum_{i,j} b_{ij} |A + sB|_{ij} \quad ,$$

From Aitken (1951, pg. 131) we have that

Pyke has shown that the matrix of L.S.T.'s of the  $P_{ij}^{1j}(t)$  is given by

$$(5.6) \quad \pi(s) = (I - q(s))^{-1} (I - h(s)) \quad .$$

Using (5.4) we obtain

$$\lim_{t \rightarrow \infty} P(t) = \lim_{s \rightarrow 0} \pi(s) = \lim_{s \rightarrow 0} \left\{ \frac{1}{s} \frac{\bar{v}'}{\bar{v}'} + \text{const.} + o(1) \right\}$$

$$\times \left\{ I - [I - \text{sdiag}(n_1, \dots, n_m) + o(s)] \right\} \\ = \frac{\bar{v}'}{\bar{v}'} \text{diag}(n_1, n_2, \dots, n_m)$$

$$= \frac{1}{\bar{v}' \bar{v}} \begin{bmatrix} v_1 n_1 & & & \\ \vdots & \ddots & & \\ v_2 n_2 & & \ddots & \\ \vdots & & & \ddots \\ v_m n_m & & & & \vdots \\ & & & & & v_m n_m \end{bmatrix} ,$$

whose  $(i, j)$ th element is

$$(5.7) \quad \lim_{t \rightarrow \infty} P_{ij}^{1j}(t) = \frac{v_j n_j}{\bar{v}' \bar{v}} \quad .$$

Thus we obtain an expression for the stationary distribution of  $P_{ij}^{1j}(t)$

in terms of the basic quantities specified by the matrix  $\bar{Q}(t)$ . Pyke

obtains the stationary distribution of  $P_{ij}^{1j}(t)$  as follows: define

$$(5.8) \quad G_{ij}^{1j}(t) = \text{Pr}\{N_{ij}^1(t) > 0\} \quad .$$

Pyke then proves that

$$(5.9) \quad \lim_{t \rightarrow \infty} P_{ij}^{1j}(t) = G_{ij}^{1j}(\infty) \frac{\gamma_{ij}^{1j}}{n_j} \quad ,$$

where  $\gamma_{ij}^{1j}$  is the mean of  $G_{ij}^{1j}$ . (5.7) seems more useful to the practitioner

since it is in terms of basic quantities rather than in terms of the mean

recurrence times.

From (5.4) we can obtain the limiting forms for the moments of several distributions discussed in previous chapters. In Section 1.4 we obtained

$$(5.10) \quad \int_0^\infty e^{-st} g_M(t) dt = m(s) = (I - q(s))^{-1} (I - I) \cdot$$

So that

$$\lim_{t \rightarrow \infty} \frac{1}{t} M(t) = \lim_{s \rightarrow 0} sm(s)$$

$$= \lim_{s \rightarrow 0} s \left\{ \frac{1}{eV'} + \text{const.} + o(1) \right\}$$

$$(5.11) \quad \frac{eV'}{V'} = \bar{V}'\bar{\eta}$$

whose  $(i, j)$  th element is

$$(5.12) \quad \lim_{t \rightarrow 0} \frac{1}{t} M_{ij}(t) = \frac{V'}{V'} \bar{V}'\bar{\eta}$$

From the definition of  $G_{ij}^1(t)$  given in (5.8) one has recursively that

$$(5.13) \quad G_{ij}^1(t) = [I - G_{jj}^1(t)] * G_{ij}^1(t) + \sum_{k=1}^m G_{ik}^1(t) * G_{kj}^1(t)$$

and

$$(5.14) \quad g_{ij}^1(s) = [I - g_{jj}^1(s)] q_{ij}^1(s) + \sum_{k=1}^m q_{ik}^1(s) g_{kj}^1(s)$$

Pyke shows that the solution to (5.14) is

$$(5.15) \quad g(s) = q(s) (I - q(s))^{-1} \{ [I - q(s)]^{-1} \}^{-1}$$

whose  $(i, j)$  th element is

Using (5.4) we can also obtain the asymptotic moments of the distributions

$$(5.18) \quad \frac{1}{V} = \frac{1}{V} \cdot \frac{1}{\bar{\eta}}$$

$$\lim_{s \rightarrow 0} \frac{\lim_{t \rightarrow \infty} M_{1j}^{jj}(t) \frac{1}{V} \frac{s}{1 + m_{1j}^{jj}(s)} + \text{const.} + o(1)}{\lim_{s \rightarrow 0} \frac{\lim_{t \rightarrow \infty} M_{1j}^{jj}(t) \frac{s}{1 + m_{1j}^{jj}(s)} + \text{const.} + o(1)}} = \frac{1}{V} \frac{s}{1 + m_{1j}^{jj}(s)} + \text{const.} + o(1)$$

in ratios of means and their limits. To this end we have which shows the equivalence of (5.7) and (5.9). One is often interested

$$\lim_{t \rightarrow \infty} P_{1j}^{jj}(t) = \frac{1}{\eta_j} = \frac{1}{V} \frac{1}{\bar{\eta}}$$

Note that if the state  $j$  is recurrent then from (5.9)

$$(5.17) \quad \frac{1}{V} = \frac{1}{V} \cdot \frac{1}{\bar{\eta}}$$

$$\frac{1}{V} = \lim_{s \rightarrow 0} s(1 + m_{1j}^{jj}(s)) \frac{1}{\eta_j}$$

Taking limits on both sides gives

$$\frac{1}{V} = \frac{s \gamma_{1j}^{jj}(s) + o(s)}{s} = \frac{s(1 + m_{1j}^{jj}(s))}{s} = \frac{1 - g_{1j}^{jj}(s)}{s}$$

So from (5.16) we have

$$(5.16) \quad g_{1j}^{jj}(s) = \frac{m_{1j}^{jj}(s)}{1 + m_{1j}^{jj}(s)}$$



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studied in Chapter II. For example,

$$\lim_{t \rightarrow \infty} \frac{1}{t} M_{\alpha\beta}^{-\alpha}(t) = \lim_{s \rightarrow 0} \lim_{s \rightarrow 0} s m_{\alpha\beta}^{-\alpha}(s)$$

$$= \lim_{s \rightarrow 0} \lim_{s \rightarrow 0} s q_{\alpha\beta}^{-\alpha}(s) \left[ \frac{1}{s} \frac{\bar{e}v'}{\bar{v}'} + \text{const.} + o(1) \right] e^{-\alpha}$$

$$= \lim_{s \rightarrow 0} p_{\alpha\beta}^{-\alpha} \frac{\bar{e}v'}{\bar{v}'}(s) \frac{\bar{v}'}{\bar{v}'} e^{-\alpha}$$

$$= p_{\alpha\beta}^{-\alpha} \frac{\bar{v}'}{\bar{v}'} e^{-\alpha}$$

$$= \frac{p_{\alpha\beta}^{-\alpha} \bar{v}'}{\bar{v}'} e^{-\alpha}$$

whose  $i^{\text{th}}$  element is

$$(5.20) \quad \lim_{t \rightarrow \infty} \frac{1}{t} M_{\alpha\beta}^{\alpha}(t) = \frac{p_{\alpha\beta}^{\alpha} \bar{v}'}{\bar{v}'} e^{-\alpha}$$

in agreement with the asymptotic expression (2.4.5) as required. We can also use (5.4) to determine limiting moments for the arbitrary interval

$(t_0, t_0 + t]$  as for example

$$\lim_{t \rightarrow \infty} \frac{1}{t} M_{jk}^i(t) = \lim_{s \rightarrow 0} \frac{s_0}{s} [m_{jk}^i(s) - m_{jk}^i(s_0)]$$

$$= \frac{1}{s_0} \lim_{s \rightarrow 0} s \left[ 1 + \frac{s_0}{s} + \left( \frac{s_0}{s} \right)^2 + \dots \right]$$

$$\times \left\{ \frac{1}{s} \frac{\bar{v}'}{\bar{v}'} + \text{const.} + o(1) - m_{jk}^i(s_0) \right\}$$

$$(5.21) \quad = \frac{1}{s_0} \frac{\bar{v}'}{\bar{v}'}.$$

Inverting with respect to  $s_0$  gives

$$(5.22) \quad \lim_{t \rightarrow \infty} \frac{1}{t} M_{jk}^i(t) = \frac{1}{\bar{v}'} \frac{\bar{v}'}{\bar{v}'} \frac{1}{t} \bar{v}'$$

which is independent of  $t_0$  and the initial state.

M.R.P.'s and Two-State M.R.P.'s.

In Chapter IV we consider the results of Chapter II for zero-order  
 is to be published in a separate paper by Wysocki and Kshirsagar (1969).  
 as in Chapter II we were concerned with intervals  $(0, t]$ . This material  
 Chapter III we treat the case of arbitrary intervals  $(t_0, t_0 + t]$  where-  
 Chapter III further generalizes the results of Chapter II. In

of this generalized cumulative process are then derived.

are made and the associated sojourn times. The distribution and moments  
 which the g-functions depend on the states between which the transitions

this distribution we consider a generalization of cumulative process in

moments of the transition count matrix  $[N_{jk}^i(t)]$  are derived. Having derived

$\bar{N}_i^j(t)$  we generalize these results in Chapter II where the distribution and

After deriving the distribution and moments of  $N_i^j(t)$  and the vector

later chapters. In this sense we have tried to make the paper self-contained.

rather to include in the first chapter those results of direct use in

no attempt to present a complete historical account of the theory, but

S.M.P.'s and M.R.P.'s number well above 100. In this paper we have made

The researcher soon realizes that the theoretical and applied papers in

for the several papers that have appeared in the literature since then.

agree that the two papers published in 1961 by Ronald Pyke were the catalyst

Those who possess a cursory knowledge of S.M.P.'s and M.R.P.'s will

## SUMMARY AND FURTHER RESEARCH

### CHAPTER VI

Chapter V utilizes a power series expansion of  $(I - q(s))^{-1}$  to derive limiting distributions and their moments in terms of the basic quantities  $Q_{ij}(t)$ .

Although several researchers have been attracted to M.R.P. theory they have severely neglected the estimation problems that arise in M.R.P.'s. The reader may consult recent bibliographies by Neuts (1968) and Fox (1968) to see that the theoretical papers discuss distribution theory almost exclusively. A recent paper by Pyke and Moore (1968) is the only paper, to this author's knowledge, that discusses the estimation problem for M.R.P.'s. Those anticipating research in this area would do well to study the results of Bartlett (1951) concerning goodness of fit tests for Markov chains and a book by Billingsley (1961) on inference in Markov Processes. The bridge between theory and application is yet to be built.

One further generalization of this paper has indications of being a fruitful avenue of research. We motivate the problem by means of a physical example from inventory models. Consider an inventory of  $n$  spare parts which supply a machine and assume that these parts have been previously used and therefore, at machine time  $t = 0$ , have ages specified by the  $n \times 1$  vector  $\bar{X}(0)$ . Assume that the distribution of failure times for the spare parts, given that the spares are  $\bar{X}(0)$  old, is not the same as the unconditional distribution; i.e., failure time distributions are not exponential. If one considers the M.R.P. associated with this spare parts problem, call it a modified M.R.P. (M.M.R.P.), then several questions come to mind.

- (1) Find the distribution and moments of  $\{Pr\{N(t) = k | \bar{X}(0) = \bar{j}(0), \bar{e}_1^T\}$  and  $\{Pr\{N(t_0, t) = k | \bar{X}(0) = \bar{j}(0), \bar{e}_1^T\}$ .

This list of research problems is presented to give the reader a flavor for the avenues of further study. Indeed one could extend the list to include a generalization of Markov Chain theory to S.M.P.'s.

(v) Instead of one machine and one spare parts inventory we have several machines each having their own spare parts inventory. What can be said about the superposition of these M.M.R.P.'s? What is the distribution of the interval between successive failures in the superposed process? If we begin observing the superposed process at  $t_0$ , then what is the distribution of time to the next renewal?

(iv) Assuming that we have a second machine of a different kind that breaks down according to failure time distribution  $G(x)$  find the distribution of  $N_t^{jk}$  in the time interval between successive failures of the second machine in the interval  $(t_0, t_0 + t]$ . We call this distribution the G-counts of the M.R.P. For  $t_0 = 0$  and  $G$  a Poisson random variable the reader is referred to Kshirsagar (1969).

$$\Pr\{N_t^I(t) = k\} , \Pr\{N_t^I(t_0, t) = k\} .$$

(iii) We may have prior probabilities for  $\bar{X}(0)$ , in which case find the distribution and moments of  $\bar{X}(0)$  ?

(ii) If  $\bar{X}(0)$  is given but we do not know how to assign an element of  $\bar{X}(0)$  to a spare part, then what can be said of the distributions

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13. ABSTRACT

Consider a particular system which at time  $t$  may be in one of a finite number of distinguishable states, labeled for convenience by  $1, 2, \dots, m$ . Once the system enters a particular state, say  $i$ , it instantaneously selects the next state to be visited, say  $j$ , with probability  $P_{ij}$ . However, transition to state  $j$  occurs after holding in state  $i$  for a random time (sojourn time) whose distribution function is  $F_{ij}(\cdot)$ . These processes are known as Semi-Markov Processes and the associated renewal process is called a Markov Renewal Process. In this paper we introduce a new counting process which at time  $t$  counts the number of times the system has made a one-step transition from state  $i$  to state  $j$ ,  $i, j = 1, 2, \dots, m$ . The matrix  $N(t)$  denotes these counts. The distribution and moments of  $N(t)$  are derived and cumulative processes associated with  $N(t)$  are discussed. We also extend these results to arbitrary intervals of the form  $(t_0, t_0 + t]$ . Certain limiting results are given and some important special cases discussed. Several open problems are given in the summary chapter.